

# **Compact Isolated High Frequency DC/DC Converters Using Self-Driven Synchronous Rectification**

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## **(ABSTRACT)**

In the early 1990's, with the boom of the Internet and the advancements in telecommunications, the demand for high-speed communications systems has reached every corner of the world in forms such as, phone exchanges, the internet servers, routers, and all other types of telecommunication systems. These communication systems demand more data computing, storage, and retrieval capabilities at higher speeds, these demands place a great strain on the power system. To lessen this strain, the existing power architecture must be optimized.

With the arrival of the age of high speed and power hungry microprocessors, the point of load converter has become a necessity. The power delivery architecture has changed from a centralized distribution box delivering an entire system's power to a distributed architecture, in which a common DC bus voltage is distributed and further converted down at the point of load. Two common distributed bus voltages are 12 V for desktop computers and 48 V for telecommunications server applications. As industry strives to design more functionality into each circuit or motherboard, the area available for the point of load converter is continually decreasing. To meet industries demands of more power in smaller sizes power supply designers must increase the converter's switching frequencies. Unfortunately, as the converter switching frequency increases the efficiency is compromised. In particular, the switching, gate drive and body diode related losses proportionally increase with the switching frequency.

This thesis introduces a loss saving self-driven method to drive the secondary side synchronous rectifiers. The loss saving self-driven method introduces two additional transformers that increase the overall footprint of the converter. Also, this thesis proposes a new magnetic integration method to eliminate the need for the two additional

gate driver magnetic cores by allowing three discrete power signals to pass through one single magnetic structure. The magnetic integration reduces the overall converter footprint.

**To my wife**  
**Elizabeth Jewel**

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# Chapter 1

## Introduction

### 1.1 Background and Objectives

From the early 1990's, with the boom of the Internet and the advancements in telecommunications, the demand for high-speed communications systems has reached every corner of the world in forms such as, phone exchanges, the internet servers, routers, and all other types of telecommunication systems. These communication systems require more data computing, storage, and retrieval capabilities at higher speeds, placing a great demand on the power system. To meet these demands, the existing power architecture must be optimized.

The old power delivery architecture was a centralized unit, which delivered the required voltages directly to the point-of-load. Unfortunately, the inherent problems of regulating multiple loads at higher currents, this approach is no longer practical. For that reason a distributed power architecture was adopted. The new architecture has a common intermediate DC bus voltage, ranging from 36- to 75-volts, nominal at 48-volts, which supports the entire system. Because each communications board requires a specific voltage lower than the bus voltage, a point-of-load (POL) converter is used. The POL converter isolates and converts the intermediate bus voltage to the required microprocessor core voltage ranging from 0.8- to 3.3-volts.

The area allocated to the POL is continuously shrinking as industry attempts to pack more functionality into each of the circuit boards. After decades of effort, the state of the art POL converters switching frequency has been pushed to several hundred

kilohertz, in an attempt to increase current within smaller footprints. Today's POL converters for the telecom industry are called a "brick". The brick type converters come in four standardized sizes, a full brick, half brick, quarter brick and an eighth brick. The majority of the POL converters have the quarter brick form factor. The size of a quarter brick is 1.45 inches X 2.28 inches.

There are many different types of isolated topologies that could be used for a compact DC/DC converter such as a quarter brick. A few examples of these topologies are a fly back, forward, push-pull, half-bridge, full-bridge etc. Today's high end telecom quarter brick power supplies need to be able to supply greater than 60A of load current at output voltages ranging from 0.8V to 3.3V. Because the secondary side conduction losses are the most significant, these topologies typically have a secondary side current doubler rectifier.

One topological candidate to supply this load is the single magnetic push-pull forward converter with a built-in input filter and coupled-inductor current doubler, proposed by Peng Xu, shown in Fig. 1.1.

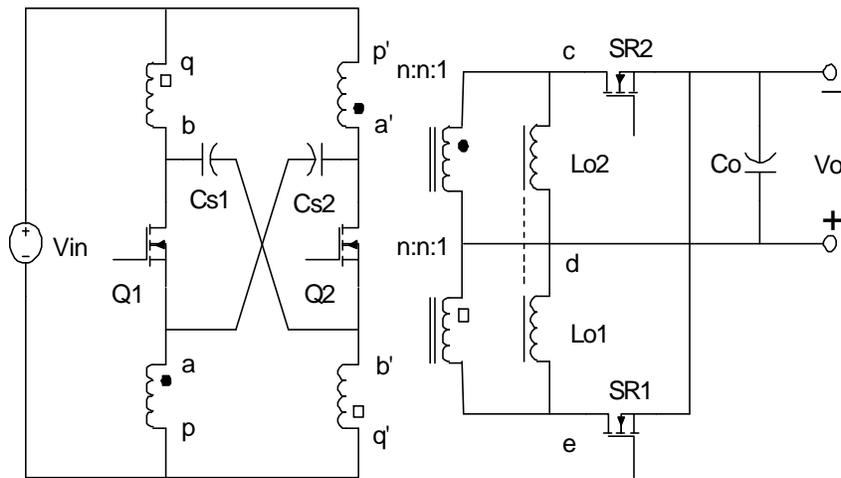


Fig. 1.1. Single magnetic push-pull forward converter with a built-in input filter and coupled-inductor current doubler

At a 100 kHz switching frequency the push-pull forward can convert 48-volts to 1.2-volts at 70 amps of load current with greater than 84 percent efficiency. The

efficiency curve for the designed push-pull forward is shown in Fig. 1.2, and a loss breakdown for the converter is shown in Fig. 1.3.

Since the switching frequency is only 100 kHz the highest losses are the winding losses and secondary side conduction losses, as seen in the loss breakdown in Fig. 1.3.

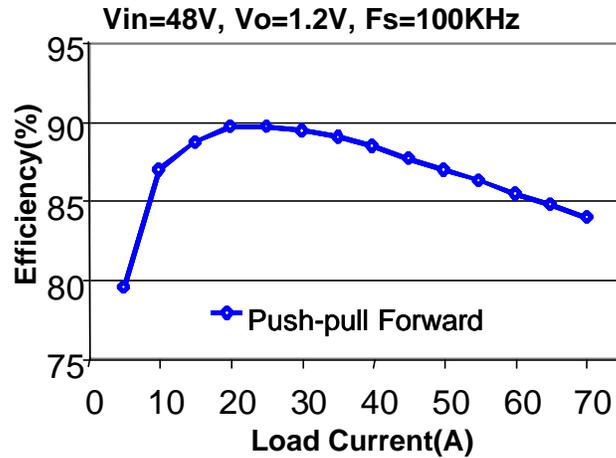


Fig. 1.2. Efficiency curve for the single magnetic push-pull forward converter with a built-in input filter and coupled-inductor current doubler

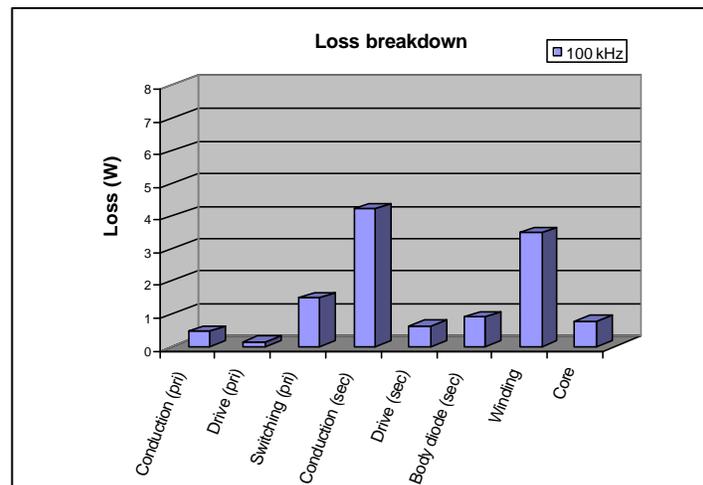


Fig. 1.3. Loss breakdown for the single magnetic push-pull forward converter with a built-in input filter and coupled-inductor current doubler operating at 1.2V / 70 A, with a switching frequency of 100 kHz.

For the push-pull forward to be a viable solution for a compact DC/DC converter, such as a quarter brick, the switching frequency must be increased to reduce the size of the passive components. Pushing the switching frequency to 500 kHz will result in an entirely different loss breakdown. A comparison between the 100 kHz and 500 kHz loss breakdowns is shown in Fig. 1.4. At 500 kHz the primary side switching losses are now the most significant. There is a five times increase in these losses because both the gate drive and body diode losses are frequency dependent. In addition the overall efficiency drops by ten percent and causes the heavy load efficiency to reduce to 74 percent for a total of 28 watts of loss.

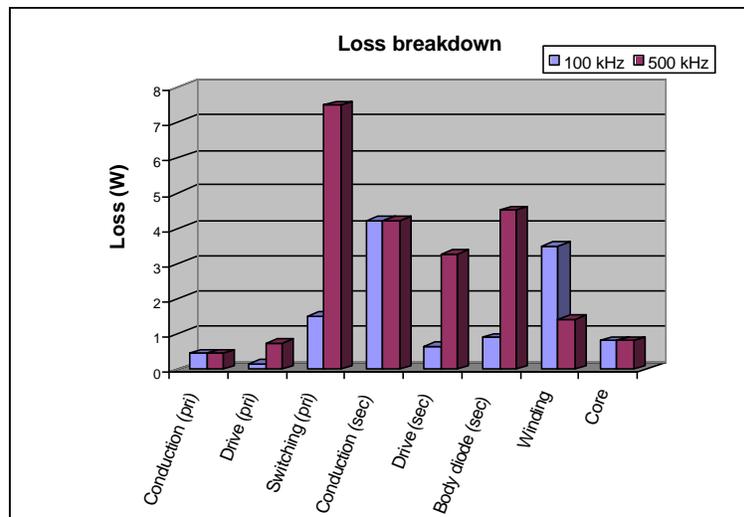


Fig. 1.4. Loss breakdown comparison for the single magnetic push-pull forward converter with a built-in input filter and coupled-inductor current doubler operating at 1.2V / 70 A, with switching frequencies of 100 kHz and 500 kHz.

Soft switching techniques such as zero-voltage-switching (ZVS) or zero-current-switching (ZCS) must be used to alleviate the highest loss bar in Fig. 1.4. The secondary side body diode losses can be saved using two methods. The first solution is to not have the body diode conduct at all or a very fast gate driver that limits the body diode conduction period. The second solution incorporates, current steering methods, or ZCS, of the secondary side synchronous rectifiers to eliminate the body diode conduction time. To save the gate driver losses, a resonant, or self-driven, gate driver can be used.

The gate driver topological selection is one of the most significant choices in a compact isolated high frequency design. The gate driver should be fast, but also have the correct drive timing that corresponds to the power delivery periods, as well as, have minimal size. Preferably, the gate driver should have loss savings properties. A conventional gate driver is considered to be the fastest, but has the highest loss and needs perfect timing. Resonant gate drivers have loss saving properties, but typically have complex designs. The state-of-the-art self-driven schemes have loss saving properties, fast switching speeds, no timing issues and simple designs. However, during the power transformer dead times in symmetrically driven topologies, the secondary side gate drive signals are disengaged causing higher body diode conduction times.

## **1.2 Thesis Outline**

This thesis is composed of four chapters.

Chapter 1 includes a brief background of the developing trends for compact isolated high frequency DC/DC converters with a focus on the telecommunications brick type form factors. This chapter also discusses the difficulties and challenges in designing secondary side gate drivers.

Chapter 2 discusses the development of a complementary-controlled full bridge converter that uses self-driven synchronous rectification.

Chapter 3 proposes the concept of integrated transformers to reduce the overall size, complexity and cost of the self-driven complementary-controlled full bridge. The integrated transformer concept is used to integrate two gate driver transformers and one power transformer onto a single EE/EI core. The leakage inductances between the secondary and primary windings are equivalent to that of three discrete transformers. This integration method is unlike other self-driven methods that couple the secondary side gate driver windings to the main power transformer. Experimental software and

hardware results are shown for both discrete transformer and integrated transformer prototypes.

Chapter 4 summarizes the conclusions and proposes ideas for future work in an extension of the integrated transformer concept beyond just integrating the two gate drive transformers with the main power transformer. The two output inductors are further integrated with the three winding transformer. Also in this chapter the light load efficiency of the self-driven complementary-controlled full bridge circuit is examined instances where a wide input voltage range is required.

## **Chapter 2**

# **Development of a Zero-Voltage-Switched Complimentary-Controlled Full Bridge with Self- Driven Synchronous Rectification**

The zero-voltage-switched (ZVS) phase-shifted full bridge is one of the most well known of the soft-switched DC/DC converters. Another, but lesser known form, of the ZVS phase shifted full bridge is the complimentary controlled full bridge. The major benefit of the complimentary controlled full bridge is the topologies ease in implementing self-driven synchronous rectification. This chapter will derive the complimentary controlled full bridge from the phase-shifted full bridge and demonstrate the enhanced design for the self-driven synchronous rectifiers.

### **2.1 Secondary Side Rectifier Topologies**

The topology selection for a compact high frequency, high current and low voltage converter is driven by thermal performance and power density. The converters efficiency must be maximized and the hot spots must be eliminated to improve the thermal performance in applications where a heat sink cannot be applied. In applications that have a large step down ratio, such as a 36 to 75 V input to less than 3.3 V output, the efficiency could be optimized with the use of a step down transformer. For applications that have an isolation transformer, the secondary side conduction losses are dominant. Synchronous rectifiers (SR's) can replace Schottky rectifiers to lower the secondary side

conduction losses. Multiple SR's can be placed in parallel to further reduce the conduction losses and to improve the secondary side thermal distribution.

Isolated topologies can have a multitude of secondary side topologies to fit any application or budget. The four most commonly seen secondary side topologies are the full-wave, forward, center-tap, and current-doubler rectifiers, as seen in Fig. 2.1.

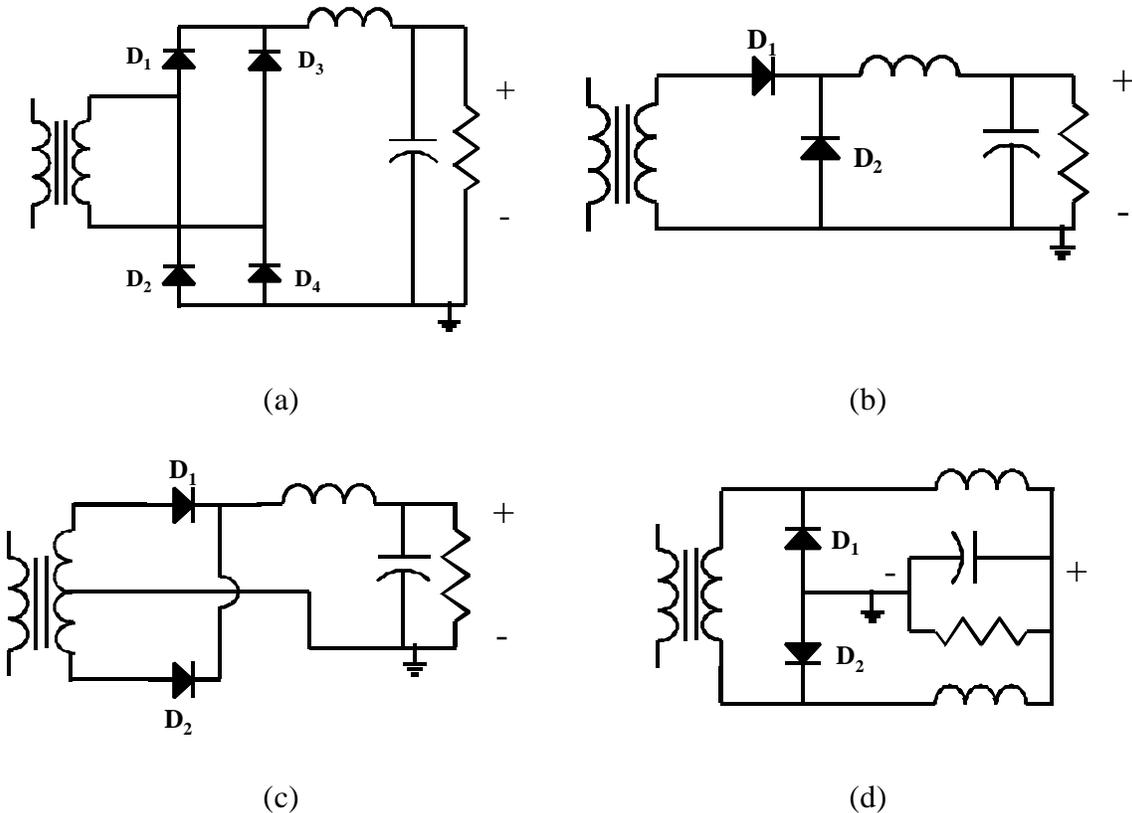


Fig. 2.1. (a) Full-wave rectifier, (b) Forward rectifier, (c) Center-tap rectifier, and (d) Current-doubler rectifier

During the power delivery period of the full-wave rectifier, Fig. 2.1a, the filter inductor current will conduct through two diodes in series. In the other three secondary side topologies, Fig. 2.1b, Fig. 2.1c and Fig. 2.1d, the full load current will conduct through a maximum of one device in series. The conduction losses of the full-wave rectifier are about twice that of the forward, center-tap and current-doubler. Although applicable to high voltage conversion, the full wave rectifier, Fig. 2.1a, is not very suitable for low voltage high current conversion or for compact, high frequency conversion.

The forward rectifier, Fig. 2.1b, has the most basic structure, and contains only a simple transformer, two switches, a single filter inductor and a capacitor. But, of the remaining three rectifier topologies, the forward rectifier is the least suitable for high current, low voltage applications because the forward rectifier has a larger filter inductor and larger rectification losses. The filter inductor is sized to support the full load current at the converter's switching frequency. During the power delivery period the inductor current will flow through  $D_1$ , and, during the freewheeling period  $D_2$  will carry the full inductor current. In essence, the losses of  $D_1$  and  $D_2$  equal that of one device carrying the full output inductor current during the entire switching period.

Compared to the forward topology, the filter inductor in the center-tap rectifier, Fig.2.1c, is designed at twice the switching frequency resulting in a smaller inductance value. The filter inductor value is inversely proportional to the effective ripple current frequency seen by the inductor, which results in a fifty percent reduction in the required inductance value. During the freewheeling period in a symmetrically driven center-tap rectifier the filter inductor current is evenly distributed between both rectifiers  $D_1$  and  $D_2$  resulting in lower rectifier conduction losses. Fig 2.2 shows an asymmetrical and a symmetrical transformer voltage waveform.

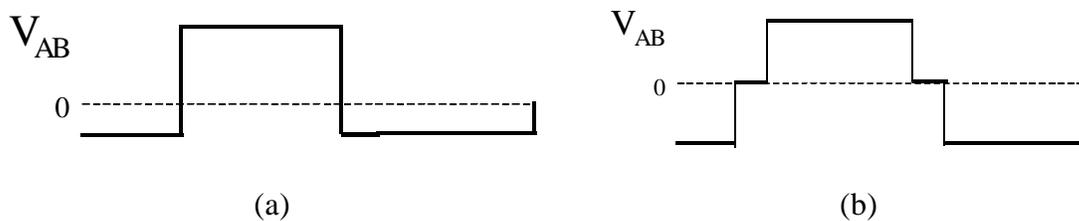


Fig. 2.2. Secondary side transformer waveforms (a) asymmetrical, (b) symmetrical

Compared to the forward and center-tap rectifiers, the current doubler has an additional filter inductor. The frequency seen by the filter inductors is the same as the switching frequency in the current doubler, but the size of these inductors can also be reduced. The filter inductors size reduction comes from a partial cancellation of the ripple current due to the interleaving properties of the current doubler. During the

freewheeling period, the load current in a symmetrically driven current-doubler will distribute evenly between each rectifier.

The current-doubler rectifier is preferred over the center-tap rectifier for use in high current low voltage applications. As opposed to the center-tap rectifier, the current-doubler has two times less current in the output inductors and secondary transformer windings. Consequently, the current-doubler has lower conduction losses through the windings of the magnetic devices. Furthermore, the output inductors and the transformer of the current-doubler topology can be integrated into a single magnetic core.

To reduce the conduction losses, MOSFET synchronous rectifiers (SR) replace the diodes illustrated in Fig. 2.1, but the circuit implementation becomes more complex. Paralleling multiple synchronous rectifiers can further reduce the conduction losses. The parallel combination of multiple synchronous rectifiers will lower the effective  $R_{ds(on)}$ , and therefore lower  $I^2 * R$  losses. In this situation, the RMS current is divided between each synchronous rectifier, the losses are distributed and the thermal burden reduced. In this vein, the gate drive losses will increase with each added synchronous rectifier.

## **2.2 External Gate Drivers vs. Self-Driven Gate Drivers**

As switching frequencies increase, the need for a low loss or energy saving gate driver circuit is necessary. For high power density designs the gate driver size should also be small. There are many different types of gate drive circuits that fall under two categories, external driven or self-driven. An external gate driver receives its signal from a control circuit, where as, a self-driven circuit receives its signal and power internally from the circuit. To demonstrate the need for a low loss gate driver, a half-bridge converter with a conventional external driven current doubler secondary was designed at 100 kHz and 1 MHz. The half-bridge converter was designed for a 48-volt input to 1.2V output at 80-amps of current. The primary side switches are Hitachi's HAT2175 and five HAT2165 synchronous rectifiers were used in parallel for each side of the current doubler. To keep the output current ripple constant at 100 kHz the output filter inductor

is 400 nH, and at 1 MHz the filter inductors are 40 nH. A loss breakdown is shown in Fig. 2.3.

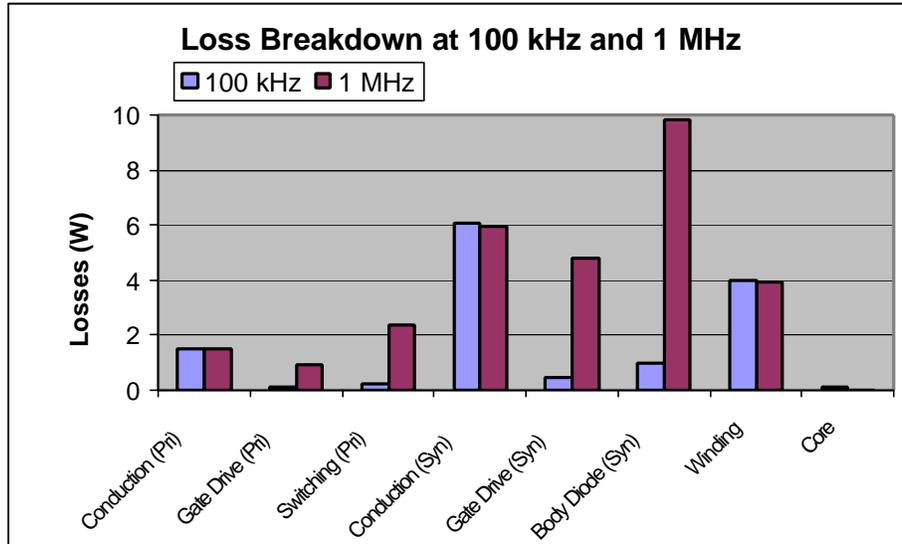


Fig. 2.3 Loss breakdown of a half-bridge converter with a current doubler secondary operated at 100 kHz and 1 MHz.

At 100 kHz the designed converter can achieve 88% efficiency at full load. After changing the switching frequency to 1 MHz there is a 10% drop in efficiency. At 1 MHz the primary side switching losses and the synchronous rectifier gate drive and body diode losses increase by ten times. The switching losses can be improved or eliminated by using soft switching techniques. The gate drive losses and body diode losses can be improved through the use of advanced gate driving techniques that save gate drive energy and minimize the body diode conduction time. A conventional gate driver is the easiest driver to implement. An example, Fig. 2.4, shows a conventional gate driver with operating waveforms.

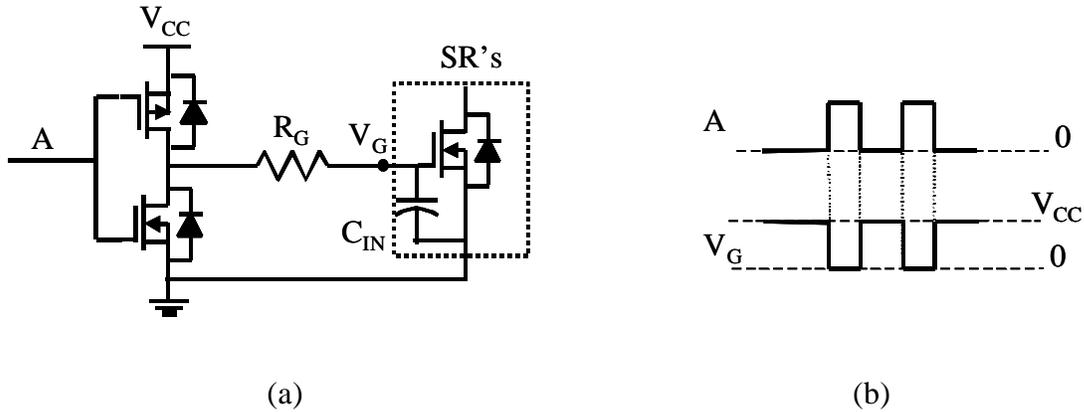


Fig. 2.4. (a) Conventional Gate Driver (b) Operating waveforms.

A conventional gate driver is considered to have the highest losses. The gate driver, in the simplest form, consists of a p- and an n-type transistor. When point A is driven to a low voltage the top p-type transistor turns on.  $V_{CC}$  then charges the gate capacitance of the synchronous rectifier through the gate resistor,  $R_G$ . When the threshold voltage of the SR is reached the SR will turn on. The gate capacitance continues to build voltage until the gate voltage reaches  $V_{CC}$ . When point A's voltage goes to a high state, the n-type transistor will conduct and the p-type transistor will open. Any existing gate charge on the SR's will be shorted to ground through  $R_G$  and the n-type transistor. The power loss associated with charging and discharging the gate capacitance of the SR's is as follows:

$$P_{LOSS\_DRIVE} = \frac{1}{2} C_G V_{CC}^2 f_S \quad (2.1)$$

Because the gate is charged and discharged during each cycle, the total power losses are twice that of (2.1). Although the conventional gate drive circuitry is very simple in design and small in size, the losses are very large and frequency dependant. As the switching frequency of the converter moves toward the MHz region, the gate drive losses will potentially negate any savings from additional paralleled synchronous rectifiers. Therefore, an alternate form of gate drive circuitry must be used.

A resonant gate driver is also an external gate driver, but has the capability to save gate drive losses. Shown in Fig. 2.5 is one example of a resonant gate driver that was proposed by Yuhui Chen.

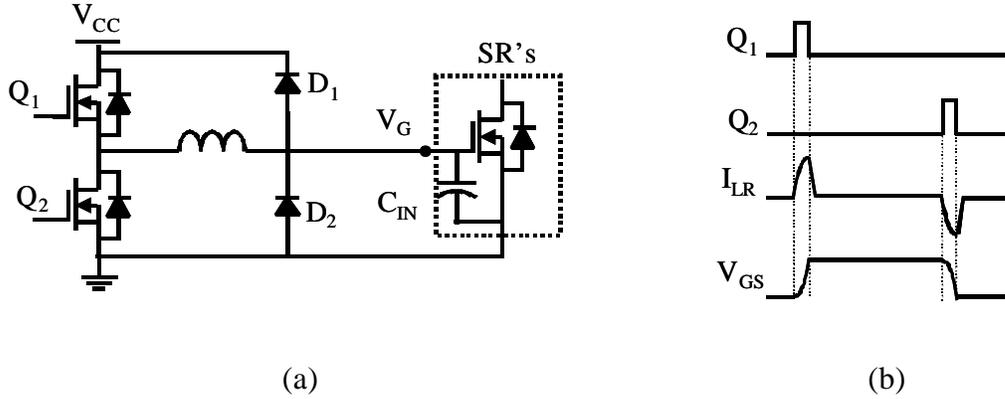


Fig. 2.5. (a) Resonant MOSFET gate driver, (b) Operating waveforms.

From Fig. 2.5, the resonant gate drive waveforms operate in a very different fashion than the conventional gate driver. During turn-on of the synchronous rectifier,  $Q_1$  is first initiated through application of  $V_{CC}$  across the resonant inductor and the gate capacitance of the SR's. The gate capacitance begins to charge in a resonant fashion up to  $V_{CC}$ . Once the gate voltage equals  $V_{CC}$  the diode  $D_1$  will begin to conduct, and clamp the gate voltage to  $V_{CC}$ . Then,  $Q_1$  will turn off. The current in the resonant inductor will continue to flow through the body diode of  $Q_2$ , and diode  $D_1$ , back to  $V_{CC}$ . This circulation recycles the gate drive energy. During the turn off of the SR's,  $Q_2$  will turn on, causing the resonant inductor and the gate capacitance to resonate. Once the gate voltage of the SR's reaches zero, diode  $D_2$  will conduct, and clamp the gate to source voltage. Then,  $Q_2$  will turn off. The current in the resonant inductor will continue to flow through  $D_2$ , and the body diode of  $Q_1$ , back to  $V_{CC}$ , thus recycling the gate drive energy. This resonant gate drive circuit can recycle the gate drive energy at both turn-on and turn-off of the synchronous rectifiers.

The resonant types of gate drive circuits have the potential for saving gate drive losses, but at a penalty of added complexity, size and difficulty in design. The described resonant gate driver needs to have multiple control signals in order to drive one set of synchronous rectifiers. In an isolated topology, these gate drive control signals need to

be derived, or computed, from the primary side and delivered to the secondary side via a transformer or an optocoupler. Also, secondary side power supply is needed.

The second form of a gate driver is a self-driven gate driver, meaning the gate drive signals are formed within the circuit itself. There are many corporations that hold patents in the area of self-driven synchronous rectification. Lucent Technologies Inc. holds some of the most commanding patents under the topic of self-driven gate drivers, and cover a very broad range of isolated topologies. The traditional self-driven gate drive scheme is to cross-couple the drive to the secondary side of the transformer.

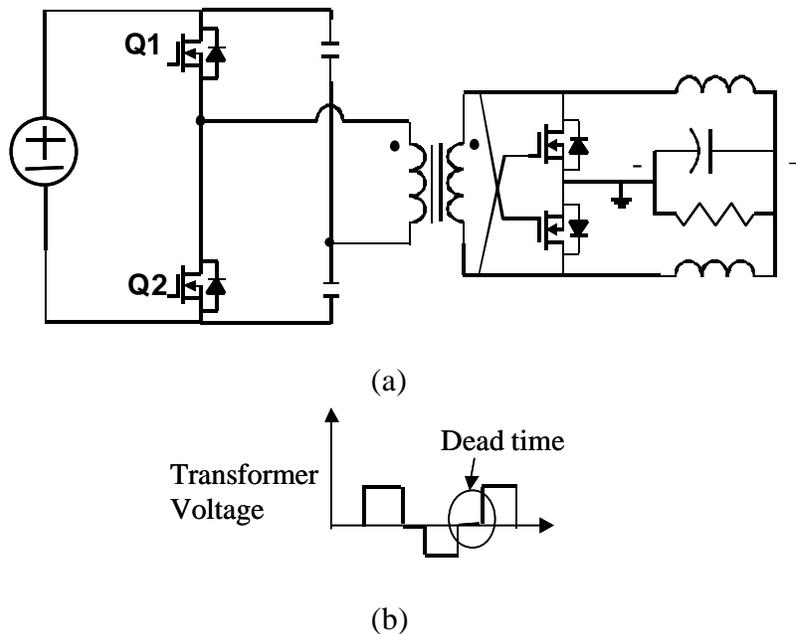


Fig. 2.6. (a) Half-bridge with self-driven current-doubler secondary, (b) Secondary side transformer waveform.

Fig. 2.6a is a half-bridge topology with a self-driven current-doubler. The transformer voltage forms the gate drive signals for the synchronous rectifiers. This gate driver is very easy to implement, but as it stands, has very little flexibility. For example, the transformer voltage commands the gate driver voltage. If the half-bridge is symmetrically driven, there will be times when zero volts are applied across the transformer leading to a dead time in the gate drive signal, as shown in Fig. 2.6b. The dead time will turn the SR's off. The output inductor currents will freewheel through the body diode, as opposed to, the channel of the synchronous rectifier increasing the losses.

If the half-bridge is asymmetrically driven, then the dead time is eliminated, but the gate drive voltage of one synchronous rectifier is larger than the other. The unbalanced gate drive voltage will cause the  $R_{ds(on)}$  of one set of synchronous rectifiers to increase, leading to higher conduction losses, and therefore, a higher operating temperature. Thus causing a lower efficiency, which results in a thermal design problem.

When the output voltage of the converter drops to a very low voltage (1.5V, 1.2V, 1V...), the secondary transformer voltage will drop to a low voltage. This drop makes driving the SR's very inefficient. Therefore, auxiliary secondary windings must be used to boost the gate drive voltage.

There are many papers and patents describing different methods of self-driven synchronous rectification that couple auxiliary windings to the transformer. For example, J.A. Cobos, proposed a self-driven level shifting method to drive the SR's on during the dead time and save gate charge losses, seen in Fig. 2.7.

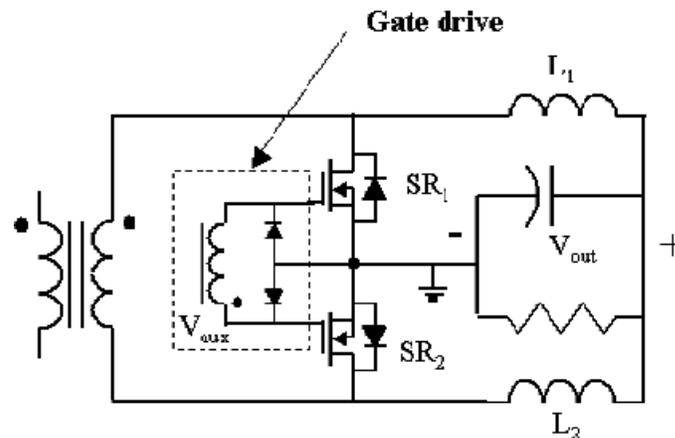


Fig. 2.7. Jose A. Cobos' self-driven circuit utilizing an auxiliary winding and two diodes.

The auxiliary winding in Fig. 2.7 is coupled with the secondary power winding along with the primary side windings. The use of the auxiliary winding allows the turns ratio between the primary and secondary power windings, and the primary and auxiliary windings, to be different. The auxiliary winding can then be optimized for the correct driving voltages. Another benefit is the level shifting properties. During the dead time,

the gate charge from one synchronous rectifier will distribute itself evenly between both SR's. During the freewheeling period both SR's will turn on and the existing gate charge will be saved. The gate drive waveforms are shown in Fig. 2.8.

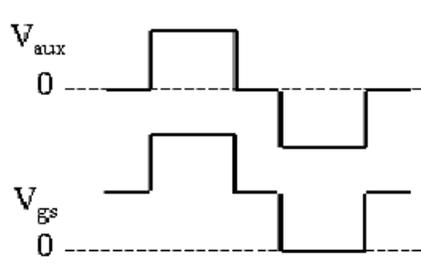


Fig. 2.8. The auxiliary winding voltage and the  $V_{GS}$  waveforms for one synchronous rectifier shown in Fig. 2.7.

In Fig. 2.8, the level shifting properties can be seen. During the freewheeling period, the gate drive voltage is one half of the voltage during the power delivery period. Even though both synchronous rectifiers turn on during freewheeling, the gate drive voltage is less than optimal, leading to higher conduction losses. Also, when the primary switch turns off, the leakage inductor of the transformer will begin to ring with the output capacitor of the primary side MOSFET. Due to the auxiliary winding being coupled directly with the power windings, the leakage energy in the power delivery path will couple into the gate driver auxiliary windings. The leakage energy will then show up at the gate of the synchronous rectifiers and cause unwanted turn-off or over voltage, ringing at the gate. At high frequency, the chances of unwanted turn-off increase. If the leakage energy or unwanted turn-off time is too large the synchronous rectifiers could be damaged, or the body diode will conduct for an excessive amount of time resulting in a lower overall efficiency. An alternate form of a self-driven gate drive circuit must be used.

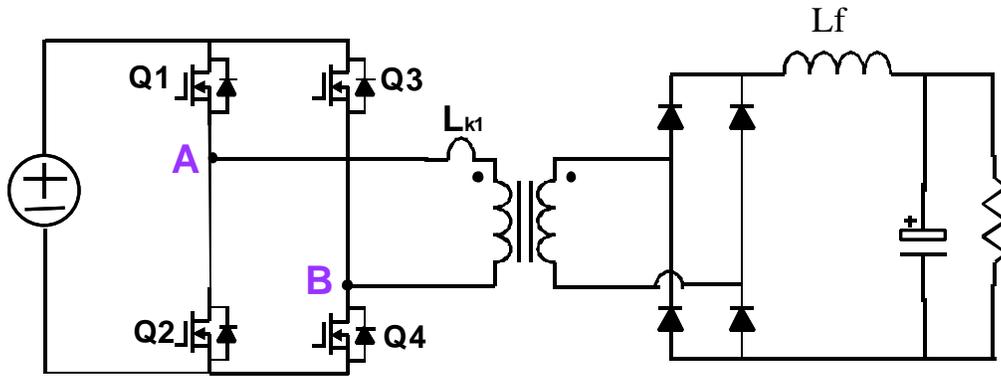
In summary, the conventional gate driver has severe losses at high frequency and with multiple synchronous rectifiers. The resonant gate driver shows promising gate driver savings, but at the price of added complexity. The self-driven gate driver shows loss savings capabilities and has a simple structure, but during dead-times has the

possibility of shutting off. One other benefit of self-driven circuits is that the gate drive signal is nearly instantaneous, thus reducing, or even eliminating, the body diode conduction time.

The externally driven gate drivers need to pass the drive signal from the primary side to the secondary side, to coordinate with the power delivery period. When the gate drive signal arrives at the secondary side there will be an additional delay time from the signal transmission. A possible fix would be to add an adaptive control scheme to compensate for the signal delay time. The adaptive control will just add further complexity to the circuit. Therefore, a simple self-driven circuit is the most advantageous gate drive circuit to use.

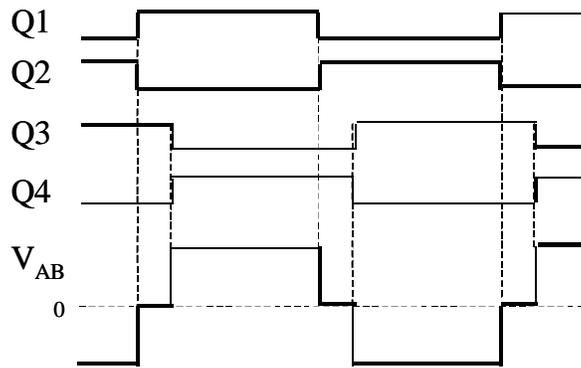
### **2.3 Zero-Voltage-Switched Phase-Shifted Full Bridge**

Referring to Fig. 2.3, the third highest frequency dependant loss at 1 MHz is the switching loss. A form of soft switching must be used to alleviate these losses. Over the years many different topologies using soft switching pulse width modulated (PWM) have been proposed. The most common soft-switched PWM topology is the phase-shifted full bridge shown in Fig. 2.9a, and the phase-shift operation is shown in Fig. 2.9b. Seen in Fig. 2.9b, switches  $Q_1$  and  $Q_2$ , as well as  $Q_3$  and  $Q_4$ , are switched in a complementary fashion with 50% duty cycle. Phase-shifting the two legs of the full bridge converter regulates the output voltage. The first leg, containing switches  $Q_1$  and  $Q_2$ , is called the active or leading leg, and as follows the second leg, containing switches  $Q_3$  and  $Q_4$  is called the passive or lagging leg. The leading leg gets its name from the rising edge of the gate signal coming before, or leading, the corresponding rising edge of the lagging leg.



(a)

Phase-Shift Operation



(b)

Fig. 2.9. (a) Phase shifted full bridge topology, (b) Ideal phase-shift operation.

The phase-shift full bridge converter is widely used because of the circuit's ability to achieve zero-voltage-switching of all four of the primary side switches. Fig. 2.10 shows the circuit timing operation for the phase-shift full bridge.

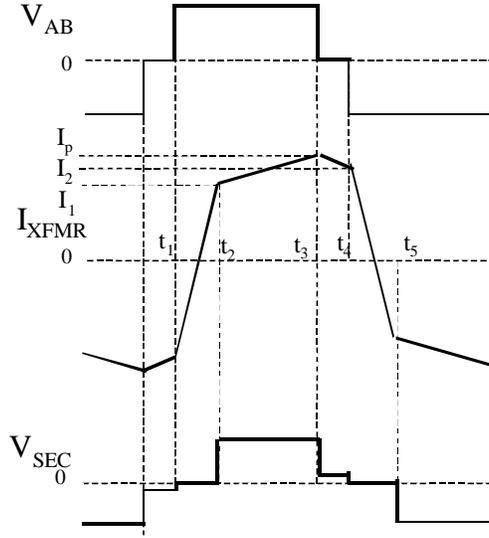


Fig. 2.10. Circuit timing operation of the phase-shift full bridge.

Starting at just before time  $t_1$ , switches  $Q_1$  and  $Q_3$  are on, and zero volts are applied across the transformer; essentially shorting out the transformer. The load current will freewheel on the secondary side with all four diodes, in Fig. 2.9a, conducting. The transformer leakage current will freewheel on the primary side, down through switch  $Q_1$ , and back up through switch  $Q_3$ .

At time  $t_1$ , switch  $Q_3$  will shut off. The leakage current in the transformer will continue to flow, and will begin to charge and discharge the output capacitor of switch  $Q_3$  and  $Q_4$ , in a sinusoidal fashion. If the leakage energy is high enough, the voltage across switch  $Q_3$  will resonate up to the bus voltage, and the voltage across  $Q_4$  will resonant down to zero. Switching  $Q_4$  at the precise moment when the voltage across the switch is zero is called zero-voltage switching (ZVS). The condition for ZVS of the lagging leg is as follows:

$$\frac{1}{2}C_{MOS}V_{IN}^2 + \frac{1}{2}C_{TR}V_{IN}^2 \leq \frac{1}{2}L_{LK}I_2^2 \quad (2.2)$$

where  $C_{MOS}$  is the sum of the equivalent output capacitances of the primary side switches  $Q_3$  and  $Q_4$ ,  $V_{IN}$  is the input voltage,  $C_{TR}$  is the transformer winding capacitance,  $L_{LK}$  is

the value of the leakage inductor and  $I_L$  is the leakage inductor current freewheeling in the primary side just after  $Q_3$  turns off.

From time  $t_2$  to  $t_3$ , switches  $Q_1$  and  $Q_4$  are on. Voltage is built up on the secondary side of the transformer and the output filter gets charged.

At time  $t_3$ , switch  $Q_1$  will turn off. The load current reflected through the transformer will begin to charge and discharge the output capacitance of  $Q_1$  and  $Q_2$ , respectively. The condition for ZVS on the leading leg is as follows:

$$\frac{1}{2}C_{MOS}V_{IN}^2 + \frac{1}{2}C_{TR}V_{IN}^2 \leq \frac{1}{2}L_{EQ}I_O^2 \quad (2.3)$$

Where  $L_{EQ}$  is the equivalent output inductor plus the leakage inductance reflected to the primary side and  $I_O$  is the reflected load current. If the load current is high enough, the voltage across  $Q_1$  will resonate up to the bus voltage, and the voltage across  $Q_2$  will resonant down to zero. When the voltage across  $Q_2$  is zero, the switch is turned on under a ZVS condition.

From time  $t_3$  to  $t_4$ , zero volts are applied across the transformer. The energy in the leakage inductor will circulate through switches  $Q_2$  and  $Q_4$ . Then at time  $t_4$ , switch  $Q_4$  will turn off. The leakage current will then resonate again with the output capacitor of  $Q_3$  and  $Q_4$ . If the condition in (2.2) is met then ZVS of switch  $Q_3$  can be realized.

Achievement of ZVS in the lagging leg (switches  $Q_3$  and  $Q_4$ ) is much more difficult than in the leading leg (switches  $Q_1$  and  $Q_2$ ) because of the smaller amount of energy in the system. The ZVS condition for the lagging leg is dependant on the leakage inductor value and the square of the leakage current. Furthermore, the ZVS condition for the leading leg depends on the sum of the leakage inductor plus the reflected output inductor and the square of the reflected load current. Therefore, ZVS of the leading leg can be effortlessly achieved.

There are two basic methods to make achieving ZVS of the lagging leg easier. The first is to reduce the output capacitor of the MOSFET switch; the second requires

enlarging the leakage inductor value in the transformer. The output capacitor of the MOSFET switch benefits the circuit by acting as a lossless snubber that will reduce the turn-off losses and limit the dv/dt of the primary side switches. Therefore, the better way to increase the ZVS range is to enlarge the effective value of the leakage inductor.

Designing a poor coupling from the primary to secondary windings, or gapping the transformer core, can increase the effective leakage of the transformer. There are two major drawbacks with using the transformer leakage as the main source of leakage inductance, (a) creating a specific and repeatable leakage inductor by designing poor coupling from primary to secondary would be very difficult if not impossible, and (b) gapping the transformer core will significantly increase the losses associated with the transformer. The preferred method for regulating the leakage inductance is to place a small series inductor with the primary side windings of the transformer. The series inductor can be optimized to achieve ZVS over almost any load range condition and the losses and inductance can be controlled.

Enlarging the leakage inductor, or placing an inductor in series with the transformer windings, can have a great benefit when it comes to achieving ZVS, but this inductance comes with a penalty: loss of duty cycle. Duty cycle loss is equal to the amount of time the freewheeling current in the resonant inductor takes to switch from time  $t_1$  to time  $t_2$ , or time  $t_4$  to  $t_5$ , in Fig. 2.10. The duty cycle loss can be calculated by the following:

$$\Delta D = \frac{\frac{n_s/n_p}{V_{IN}} \frac{T}{L_{LK}} [2I_o - \frac{V_o}{L_o} (1-D) * \frac{T}{2}] \quad . \quad (2.4)$$

Where T is the period and  $\Delta D$  is the duty cycle loss. Duty cycle loss is detrimental to circuits working at high frequency because the duty cycle loss is nearly fixed. If the duty cycle loss is held constant, and the switching frequency increases, the available time period for power delivery decreases by the amount of duty cycle loss. Therefore, when a large duty cycle is needed, regulating the output voltage can become a problem.

The time required to achieve ZVS is contained in the duty cycle loss. The period for ZVS is easily calculated as one fourth of the resonant time of the equivalent series inductance and the total equivalent parallel capacitance of the MOSFET switch as follows:

$$dt = \left(\frac{1}{4}\right) * \left(\frac{1}{2p\sqrt{LC}}\right)^{-1} \quad (2.5)$$

Due to the different equivalent series inductors for the lagging and leading legs, the required time for ZVS operation is different for each leg. For the lagging leg, the total series inductance is the equivalent leakage inductance. In the leading leg, the total series inductor is the sum of the leakage inductance and the reflected load inductor. Therefore, the time required for ZVS in the leading leg is much shorter than that of the lagging leg. A dead time, corresponding to the required time for ZVS operation, is placed between the turn-off of one switch and turn-on of the complementary switch.

Equations (2.2) and (2.3) show the condition for ZVS as an inequality. It does not mean that there is no benefit to the phase shift full bridge if these inequalities are not met. Partial ZVS can be achieved, as seen in Fig 2.11. Also, there is no added benefit for exceeding the inequalities. Greatly exceeding the inequalities in (2.2) and (2.3) will lead to an increase in the circulating energy in the system, which, in turn, increases the conduction losses. Therefore, a trade-off between switching losses and conduction losses must be made. The load condition that ZVS is achieved must be decided.

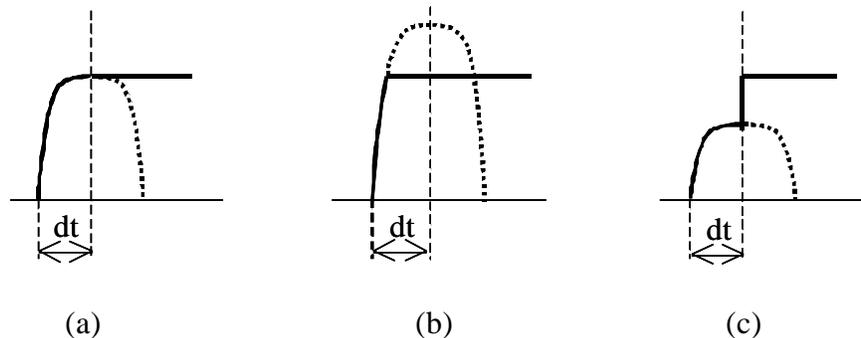


Fig. 2.11. (a) Full ZVS operation, (b) ZVS met, extra circulating energy (c) Partial ZVS, not enough energy to meet ZVS requirement.

## 2.4 Complementary-Controlled Full Bridge with a Current-Doubler

At higher switching frequencies the phase-shifted full bridge is an excellent choice because of its soft-switching capabilities. At higher output currents and lower output voltages, the current doubler rectifier with synchronous rectification is superior, though using the combination of the phase-shift full bridge and the current doubler rectifier does not lead to an easy implementation of self-driven synchronous rectification. The desire to use a self-driven scheme is previously explained. However, with a modification in the control strategy of the phase-shifted full bridge, the lesser-known complementary-controlled full bridge can be derived, as seen in Fig. 2.12.

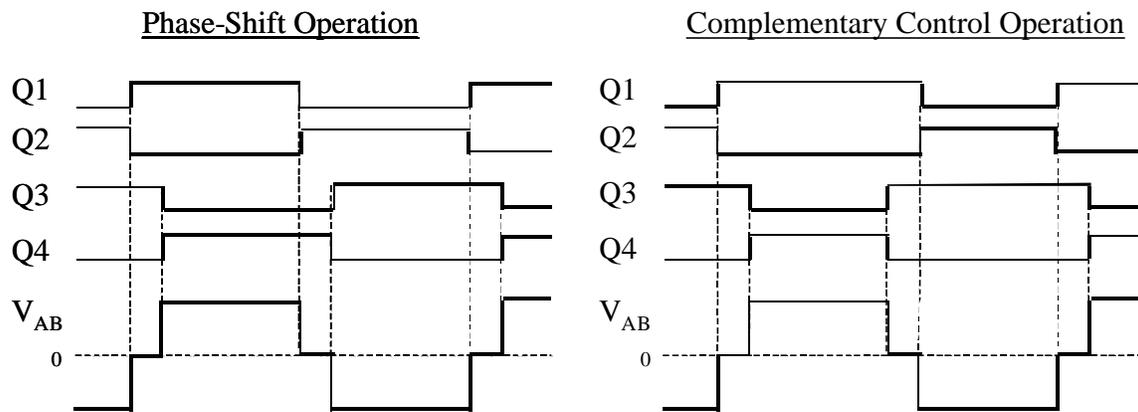


Fig. 2.12. Comparison of the phase-shift and the complementary control operations.

The operation of the complementary-controlled full bridge is nearly the same as the phase shift controlled full bridge, except that the duty cycle is formed by a controlled switch rather than phase-shifting the two legs of the converter. The complementary-control full bridge with the current doubler secondary is depicted in Fig. 2.13. From the complementary control diagram in Fig. 2.12 and Fig. 2.13, switches  $Q_2$  and  $Q_4$  are the duty cycle control switches and  $Q_1$  and  $Q_3$  are the freewheeling switches.

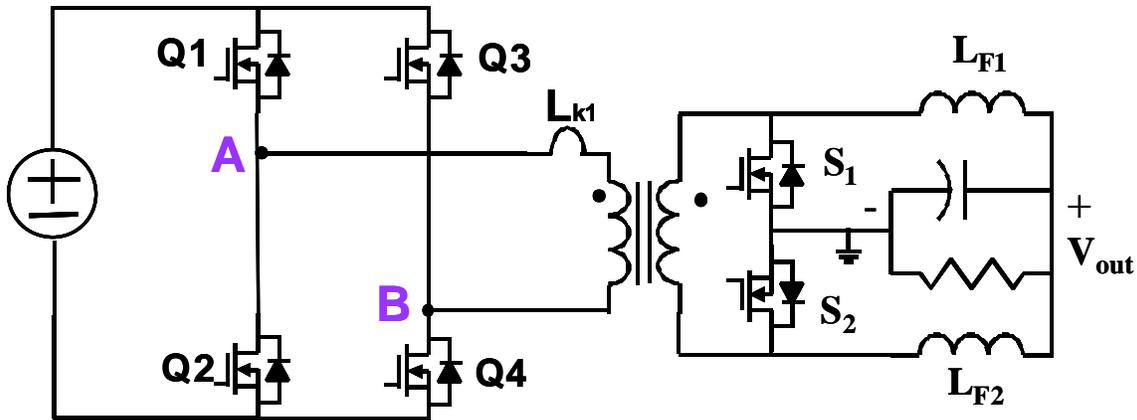


Fig. 2.13. The complementary-control full bridge with a current-doubler secondary side.

The freewheeling switches  $Q_1$  and  $Q_3$  are both turned on for greater than one half of the switching period. If the situation was reversed and  $Q_1$  and  $Q_3$  were the control switches while  $Q_2$  and  $Q_4$  were the freewheeling switches, the complementary-controlled operation would be preserved. For the duration of this work,  $Q_2$  and  $Q_4$  will be referred to as the control switches and  $Q_1$  and  $Q_3$  as the freewheeling switches.

The full circuit operation of the complementary-control full bridge with the current doubler secondary can be seen in Fig. 2.14.

Just before time  $t_b$ , the primary side of the circuit has switch  $Q_1$  on and all other primary switches are off. On the secondary side both synchronous rectifiers,  $S_1$  and  $S_2$ , are also on, however,  $S_1$  is beginning to turn off. The output inductor current in  $L_{F1}$  is freewheeling through  $S_1$  and  $L_{F2}$  is freewheeling through  $S_2$ . The current in the leakage inductor  $L_{K1}$  is freewheeling through both  $S_1$  and  $S_2$ , and then returns through the transformer. Meanwhile, the leakage current is charging the output capacitance of  $Q_3$  and discharging that of  $Q_4$ .

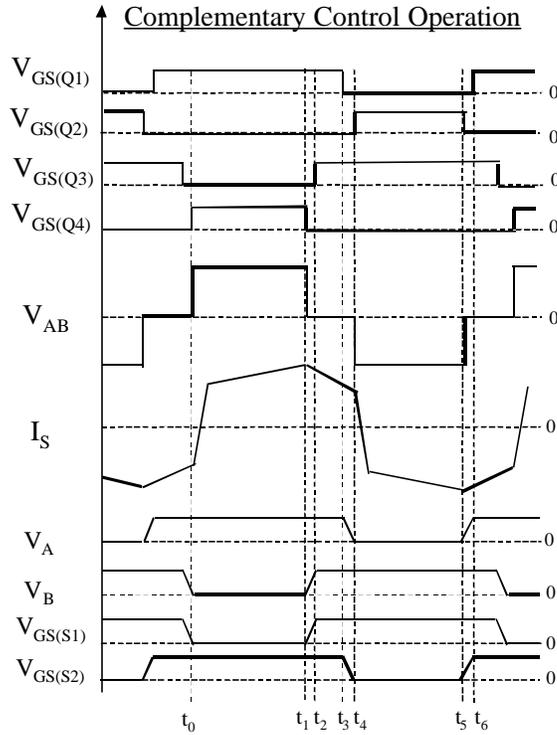


Fig. 2.14. Operation of the complementary-controlled full bridge with the current-doubler secondary side.

At time  $t_0$ , the switch  $Q_4$  turns on and  $S_1$  is turned off. If the energy stored in the leakage inductor, at time  $t_0$ , is large enough to discharge the output capacitance of  $Q_4$  to zero, then ZVS operation is achieved. The condition for ZVS of switch  $Q_4$  is identical to the condition for ZVS for the lagging leg in the phase-shift full bridge, given in (2.2). At the turn on of  $Q_4$ , the bus voltage is placed across the transformer and the leakage inductor. Once the leakage inductor current changes direction, power is delivered to the secondary side and  $L_{F1}$  begins to charge. The power delivery will continue until time  $t_1$  and  $S_2$  conducts the full load current. One half of the load current is delivered from the primary side and the other half is from the discharge of the filter inductor,  $L_{F2}$ .

At time  $t_1$ , the power delivery period is over;  $Q_4$  begins to turn off and  $S_1$  begins to turn back on. The current in the leakage inductor at this point is equivalent to the peak inductor ripple current in  $L_{F1}$ . The current in  $L_{F2}$  will continue to freewheel through  $S_2$ . Once  $Q_4$  turns off, the reflected filter inductor current in  $L_{F1}$  begins to charge the output capacitance of  $Q_4$  and discharge that of  $Q_3$ . If the circulating energy is large enough to

charge the output capacitor of  $Q_4$  to the bus voltage, then ZVS of  $Q_3$  is achieved. The condition for ZVS of  $Q_3$  is equivalent to that of the leading leg in the phase-shift full bridge, given in (2.3). Because the secondary stage is the current doubler the load current,  $I_0$  represented in (2.3) is only one half of the full load current.

From time  $t_1$  to  $t_2$ , the resonant action between the output capacitors of switches  $Q_3$  and  $Q_4$  and the reflected filter inductor takes place. At time  $t_2$ , switch  $Q_3$  and  $S_1$  turn on and the freewheeling period begins. The current in the leakage inductor freewheels around the primary side through switch  $Q_1$ , through the transformer and back through switch  $Q_3$ . The load current will freewheel through  $S_2$ , and the current in the leakage inductor  $L_{K1}$  will freewheel through both  $S_1$  and  $S_2$  and return through the transformer.

At time  $t_3$ , switch  $Q_1$  will turn off. The leakage inductor will begin to charge and discharge the output capacitors of  $Q_1$  and  $Q_2$ , respectfully.

At time  $t_4$ , switch  $Q_2$  will turn on and  $S_2$  will turn off. Again, if the leakage inductor energy is adequate, then ZVS of  $Q_2$  can be achieved. The conditions for ZVS of  $Q_2$  is the equivalent to that of  $Q_4$  and (2.2). Because  $Q_3$  and  $Q_2$  are both on, the bus voltage is applied across the transformer and the leakage inductor. Once the leakage current changes direction, power is supplied to the secondary side until time  $t_5$ . One half of the load current is delivered from the primary side charging  $L_{F2}$  and the other half is from the discharge of the filter inductor  $L_{F1}$ . Switch  $S_1$  will carry the full load current.

Right at time  $t_5$ , switch  $Q_2$  will turn off and  $S_2$  will begin to turn on. The output capacitor of  $Q_2$  and  $Q_1$  will charge and discharge, respectfully, by the reflected current in  $L_{F2}$ . A resonant action will occur between the output capacitor of  $Q_1$ ,  $Q_2$  and the reflected output inductor.

At time  $t_6$ , the switch  $Q_1$  will turn on. If the circulating energy is large enough to charge the output capacitor of  $Q_2$  to the bus voltage, then ZVS of  $Q_1$  is obtained. The condition for ZVS of  $Q_1$  is the same as that for  $Q_3$ .

From the previous analysis, the top switches  $Q_1$  and  $Q_3$  correspond to the leading leg in the phase-shift full bridge, while  $Q_2$  and  $Q_4$  correspond to the lagging leg. As

follows, ZVS can be achieved for each of the four switches. However, ZVS is harder to achieve in the control switches. The methods reported in section 2.3 for improving the ZVS range, as well as, all other prior methods for extending the ZVS range can be applied to the complementary-controlled full bridge.

Although the complementary-controlled full bridge has the same basic characteristics at the phase-shift full bridge, there is one added benefit that is an easy implementation for self-driven synchronous rectification.

From Fig. 2.14, Dr. Ming Xu observed that the gate drive signals for the synchronous rectifiers track the voltages at point A and point B in Fig. 2.13. Point A and point B correspond to the drain to source voltage of the control switches  $Q_2$  and  $Q_4$ , respectively. The voltage at point A has the identical waveform as the gate signal,  $V_{GS}$ , of  $S_2$ , and the voltage at point B has the identical waveform as the gate signal,  $V_{GS}$ , of  $S_1$ . Therefore, the following self-driven scheme, seen in Fig. 2.15 was derived.

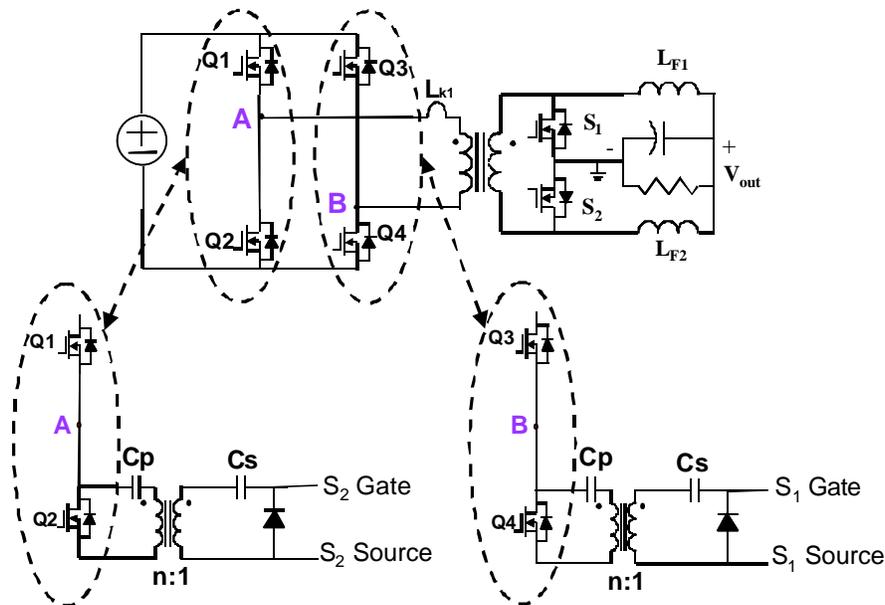


Fig. 2.15. Complementary-Controlled Full Bridge with Self-Driven Synchronous Rectifiers.

The self-driven circuit shown in Fig. 2.15 has a very simple structure with each leg consisting of two capacitors, a transformer and a diode. Because the gate driver is not

coupled to the power transformer, the synchronous rectifiers are able to track the primary side voltages at point A and B very quickly and the signals have a very clean waveform.

## 2.5 Steady-State Operation of the Proposed Self-Driven Scheme

The preceding sections presented the different secondary side rectifier topologies, the differences between external and self-driven synchronous rectifier gate drivers, and derived a new method for self-driven synchronous rectification. In this section, the operation of the proposed self-driven scheme is presented.

As discussed previously, at high frequencies the gate driver and body diode losses of the synchronous rectifiers combine to be the most significant portion of the overall losses. However, improving the gate drive speed and timing, as well as, utilizing loss saving gate drive techniques can significantly reduce these losses. The proposed self-driven circuit can save gate-driving losses and significantly reduce the body diode conduction time and losses.

To facilitate the basic operation of the self-driven circuit, Fig. 2.16 shows the path of the gate drive current during a completely non-ZVS condition for the turn-on and turn-off of the synchronous rectifiers.

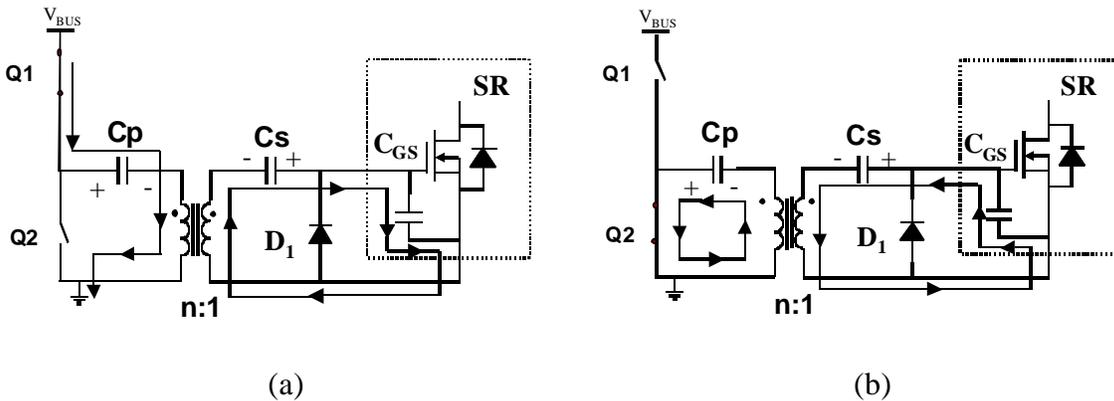


Fig. 2.16. Path of the gate drive current during complete non-ZVS (a) turn on and (b) turn off of the synchronous rectifiers.

Starting from the initial turn on of the circuit, all capacitor voltages are zero. The voltage of the primary side capacitor  $C_p$  is a function of the duty cycle and the input voltage as follows:

$$V_{CP} = V_{IN} * (1 - D) \quad . \quad (2.6)$$

The secondary side capacitor  $C_s$  voltage is a function of the primary side capacitor voltage of  $C_p$  and the turns ratio as follows:

$$V_{CS} = \frac{V_{CP}}{n} - V_D \quad . \quad (2.7)$$

Where  $V_D$  is the forward voltage drop of the diode  $D_1$ .

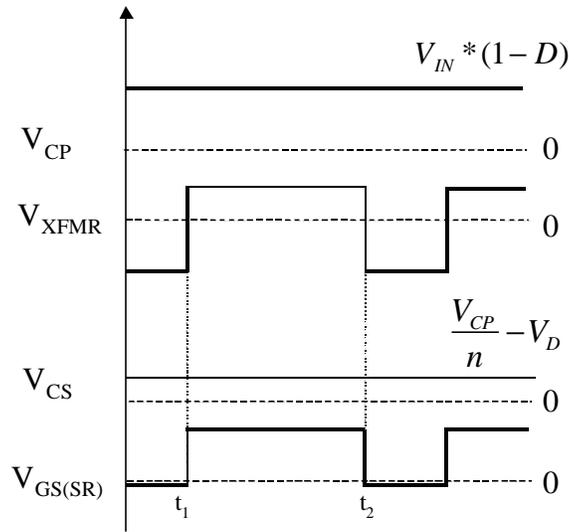


Fig. 2.17. Operating waveforms for the self-driven circuit under a non-ZVS operating condition.

For a non-ZVS condition of the primary switches, the turn-on of the synchronous rectifiers begins when the top switch  $Q_1$  turns on at time  $t_1$ . The operating waveforms are shown in Fig. 2.17. The bus voltage is then applied across the capacitor  $C_p$  and the primary windings of the gate drive transformer. The voltage across the primary side gate drive windings is equal to the bus voltage minus the voltage of  $C_p$ . On the secondary side

of the gate drive transformer, the gate voltage of the synchronous rectifiers is equal to the secondary gate drive transformer voltage added with the series voltage of  $C_S$ .

The turn-off of the synchronous rectifiers is initiated with the turn-on switch  $Q_2$  at time  $t_2$ . The capacitor voltage of  $C_P$  is then applied across primary winding. The voltage of  $C_P$  is reflected to the secondary side of the transformer. Therefore, the gate voltage of the synchronous rectifiers will equals  $V_{CP} / n - V_{CS}$ , which results in a negative voltage equal to the forward voltage drop of  $D_1$  being applied to the gate of the synchronous rectifiers. Therefore, the synchronous rectifiers will turn-off. During a non-ZVS process of the primary switches all of the synchronous rectifier gate charge is lost. However, when the self-driven scheme is used in conjunction with a ZVS mode of operation the gate drive losses can be saved.

The key ZVS operating waveforms of the complementary-controlled full bridge with self-driven synchronous rectifiers are shown in Fig. 2.18 and the topological stages are shown in Fig 2.19 through Fig. 2.24.

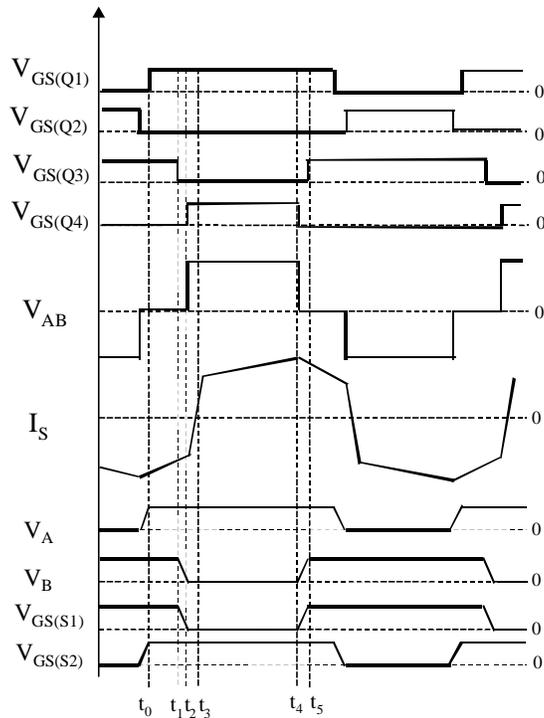


Fig. 2.18. Key operating waveforms for the self-driven complementary-controlled full bridge detailing one switching cycle of the self-driven circuit.

The converter is in a freewheeling state just before time  $t_0$ , Fig. 2.18. At time  $t_0$ ,  $Q_1$  turns on. Both of the output synchronous rectifiers and the primary side top switches  $Q_1$  and  $Q_3$  are now on. The energy stored in the transformer leakage inductor and/or primary side series inductor is freewheeling, and is represented as  $L_{EQ}$  in Fig. 2.19. The total energy in  $L_{EQ}$ , just before time  $t_0$ , is equivalent to the current in the filter inductor,  $L_{F2}$ . At this time the current in  $L_{F2}$  is as follows:

$$i_{LF2} = \frac{I_o}{2} + \frac{(1-D)}{2 * F_s} \quad (2.8)$$

Although both synchronous rectifiers are on,  $S_1$  will carry the entire load current consisting of  $I_2$  and  $I_3$  in Fig. 2.19 and  $S_2$  will carry no current. The voltage across the primary side of the transformer is zero, essentially a short circuit. Just after time  $t_0$ , until  $t_1$ , the current in the leakage inductor begins to linearly decrease causing a difference in current shared between the leakage inductor and  $L_{F2}$ .

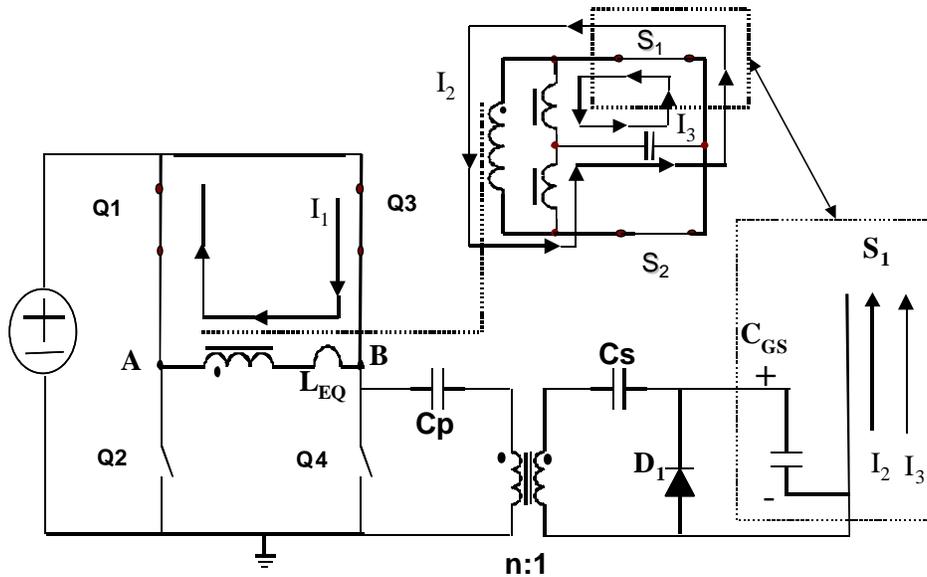


Fig. 2.19. Topological stage for the self-driven circuit at time  $t_0$ .

At time  $t_1$ , in Fig. 2.18 and 2.20, switch  $Q_3$  begins to open. The leakage energy will begin to charge the output capacitor  $C_3$  of  $Q_3$  and discharge the output capacitor  $C_4$

of  $Q_4$ . As a result, the voltage at point B, which corresponds to the drain-to-source voltage of  $Q_4$ , is discharged to zero. The leakage inductor  $L_{EQ}$  will begin to turn off the synchronous rectifier  $S_1$ , and if ZVS occurs, the gate charge of  $S_1$  is recovered. From a circuit point, of view the input capacitor of switch  $S_1$  is effectively in parallel with  $C_4$ , acting as a lossless snubber. The current in the leakage inductor during ZVS is as follows:

$$i_{LEQ} = I_1 = I_{1a} + I_{1b} + I_{5b} \quad (2.9)$$

Where the currents  $I_1$ ,  $I_{1a}$ ,  $I_{1b}$ , and  $I_{5a}$  are depicted in Fig. 2.20.

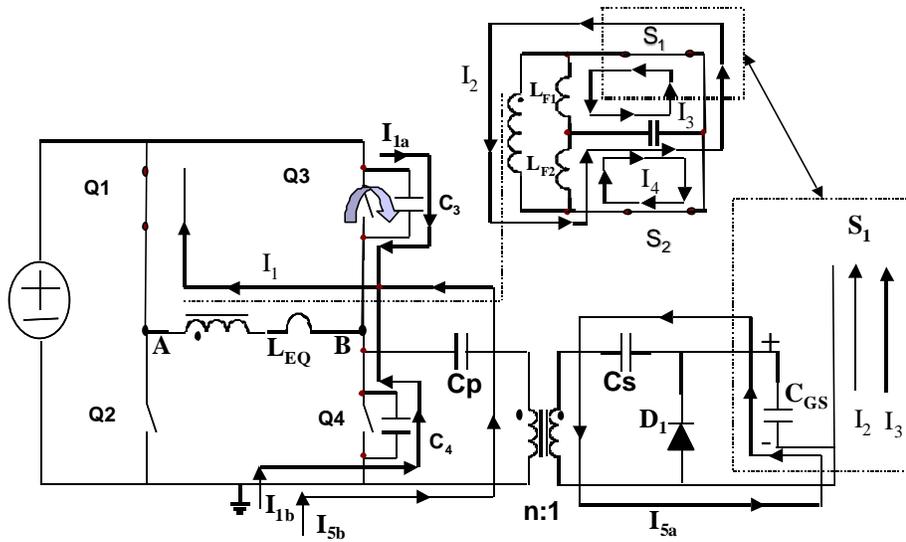


Fig. 2.20. Topological stage for the self-driven circuit at time  $t_1$ .

From time  $t_1$  to  $t_2$ , the resonant action between the leakage inductor and the parallel combination of  $C_3$ ,  $C_4$  and  $C_{gs(S1)}$  occurs. The time  $t_2$  occurs at one fourth of the resonant period. The energy needed to obtain ZVS is as follows:

$$E_{ZVS} = \frac{1}{2} (C_{MOS} + \frac{C_{GS}}{n^2}) V_{IN}^2 \quad (2.10)$$

If the energy given in (2.10) is exceeded, before time  $t_2$ , then the body diode of  $Q_4$  will begin to conduct and essentially zero volts are across  $Q_4$ . Then ZVS can be obtained. However, during the body diode conduction time of  $Q_4$ , the synchronous

rectifier  $S_1$  is turned off, allowing the body diode of  $S_1$  also to conduct. Having the body diode of  $S_1$  conduct is detrimental to the overall efficiency of the circuit. Therefore, the body diode conduction time of  $Q_4$  and  $S_1$  must be minimized.

At time  $t_2$ ,  $Q_4$  is closed, as seen in Fig. 2.21. If the ZVS operation was met then  $S_1$  was already turned off. However, if ZVS was not met, then switch  $S_1$  will begin to turn off. Assuming a non-ZVS situation, once  $Q_4$  turns off, the voltage of  $C_p$  will apply across the gate drive transformer. Because the voltage of  $C_s$  is a function of  $C_p$  and the turns ratio, given in (2.6) and (2.7), then the voltage applied to  $C_{GS}$  will be  $D_1$ . To have fast conduction of  $D_1$  a Schottky diode is used. The typical forward drop of the Schottky diode is 0.3 V. Therefore,  $S_1$  will turn off.

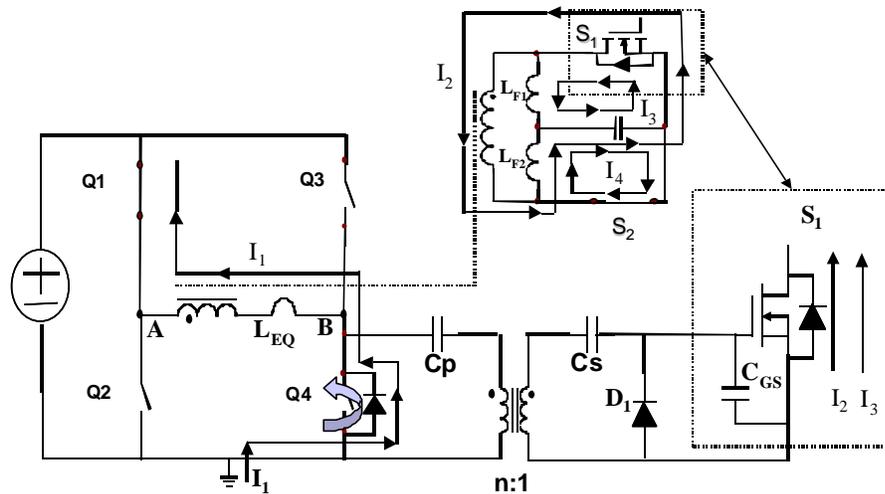


Fig. 2.21. Topological stage for the self-driven circuit at time  $t_2$ .

From time  $t_2$  to  $t_3$ , the bus voltage is applied across the primary side of the transformer. The current in the leakage inductor is forced to change directions. Once the leakage inductor changes direction, power is delivered to the secondary side at time  $t_3$ . The inductor  $L_{F1}$  then begins to charge and  $S_2$  conducts the full load current. If during  $t_2$  to  $t_3$ , switch  $S_1$  is not fully turned off, a shoot through current can occur. The timing issues and occurrences of the shoot through current are discussed in the following section.

At time  $t_3$ , the voltage of  $C_S$  will track  $C_P$  according to (2.7). This current is shown in Fig. 2.22 as  $I_{3a}$  and  $I_{3b}$ .

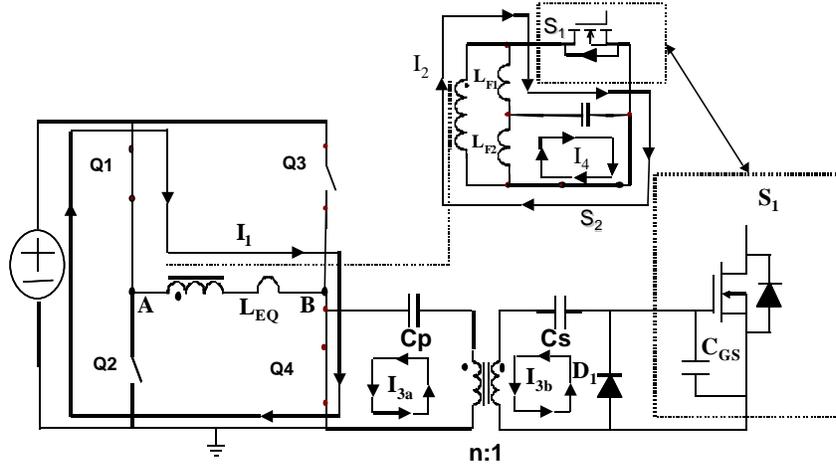


Fig. 2.22. Topological stage for the self-driven circuit at time  $t_3$ .

At time  $t_4$ , the power delivery period ends by  $Q_4$  turning off. From time  $t_4$  to  $t_5$ , the reflected current in  $L_{F1}$ , represented as  $I_1$  in Fig. 2.23, charges and discharges the capacitors  $C_4$  and  $C_3$ , respectively. Point B follows the drain-to-source voltage of  $Q_4$ . As Point B resonates up to the bus voltage, a positive voltage is applied across the gate drive transformer. Then  $I_1$  will also begin to charge the gate of  $S_1$ .

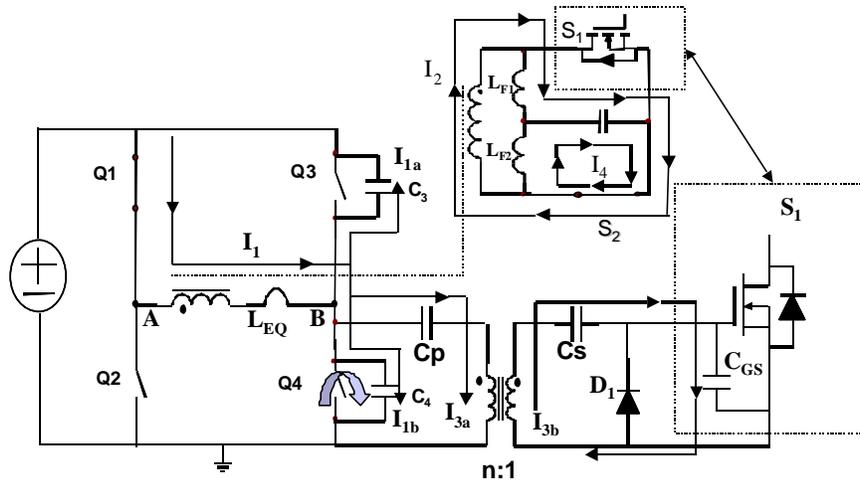


Fig. 2.23. Topological stage for the self-driven circuit at time  $t_4$ .

Time  $t_5$  is tuned to occur precisely at one fourth of the resonant time of the reflected filter inductor  $L_{F1}$ , the parallel combination of  $C_3$ ,  $C_4$ , and the reflected input capacitance of  $S_1$ . If ZVS is obtained, then the circulating current in the circuit turns on  $S_1$ . The required energy for ZVS is given in (2.10).

Under a non-ZVS condition at time  $t_5$  the voltage at point B is raised to the bus voltage by means of  $Q_3$ . Therefore, the bus voltage less the voltage of  $C_p$  is applied across the gate drive transformer. On the secondary side, the series voltage of  $C_5$  and the secondary gate drive voltage are placed across the gate to source of  $S_1$ . Then,  $S_1$  is turned on.

Once  $Q_3$  is turned on, the circuit begins a freewheeling state with  $Q_1$ ,  $Q_3$  and both synchronous rectifiers on. At time  $t_5$ , shown in Fig. 2.24, the leakage inductor current  $I_l$  freewheels on the primary side, and on the secondary side,  $S_2$  carries the full load current. After time  $t_5$ , the leakage current begins to dissipate and the switch  $S_1$  begins to carry the difference between the filter inductor current in  $L_{F1}$  and the leakage current.

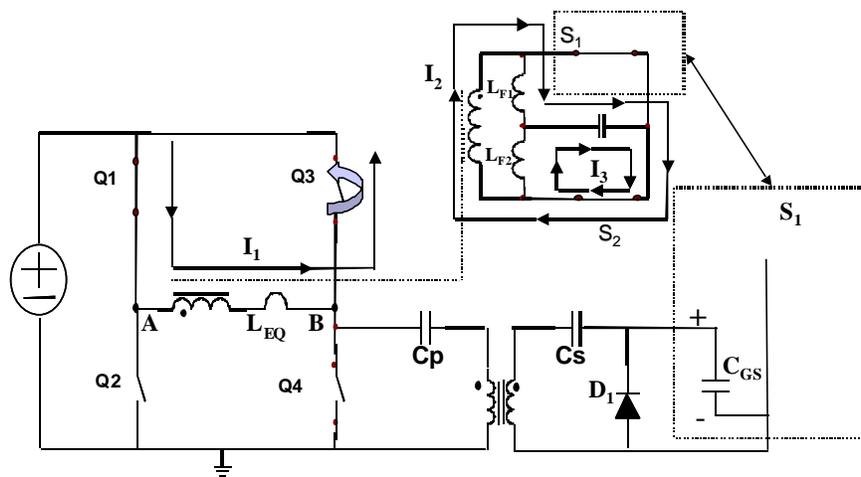


Fig. 2.24. Topological stage for the self-driven circuit at time  $t_5$ .

## 2.6 Transient Operation of the Self-Driven Scheme

The preceding section presented the operation of the self-driven scheme neglecting the transient operation and the effects of the parasitic components in the self-driven circuit. In this section, the parasitic components are taken into account in a further examination of the self-driven circuit.

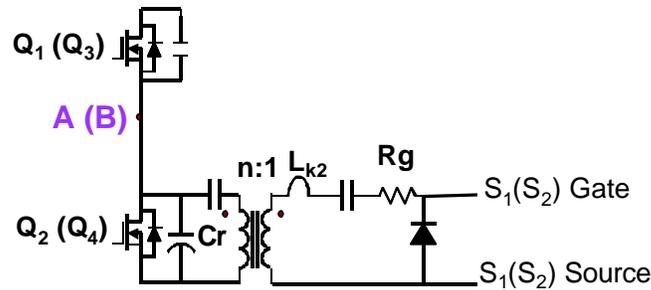


Fig. 2.25. Self-driven gate drive circuit with damping resistor.

The addition of the transformer leakage and parasitic trace inductances causes an oscillation between the leakage inductors and the gate capacitor of the synchronous rectifiers. An additional gate resistor  $R_g$  is used to dampen these oscillations, seen in Fig. 2.25.

During a ZVS turn on of a top switch, the reflected load current will charge the equivalent primary side snubber capacitor and the gate capacitor of the synchronous rectifier. During a ZVS turn on of a bottom switch, the leakage inductor of the main transformer will discharge the equivalent primary side snubber capacitor and the gate capacitor of the synchronous rectifier. The equivalent self-driven circuit diagram during turn-on and turn-off of the synchronous rectifiers is shown in Fig. 2.26.

From the circuit diagrams, it is shown that the self-driven circuit acts as a current-source driver during turn-on and a resonant type during turn-off. In actuality, during turn-on the current source in Fig. 2.26a is the reflected filter inductor, which has very high impedance, and therefore, can be modeled as a constant current source. The

damping gate resistor and leakage inductor  $L_{k2}$  directly affects the amount gate driver savings.

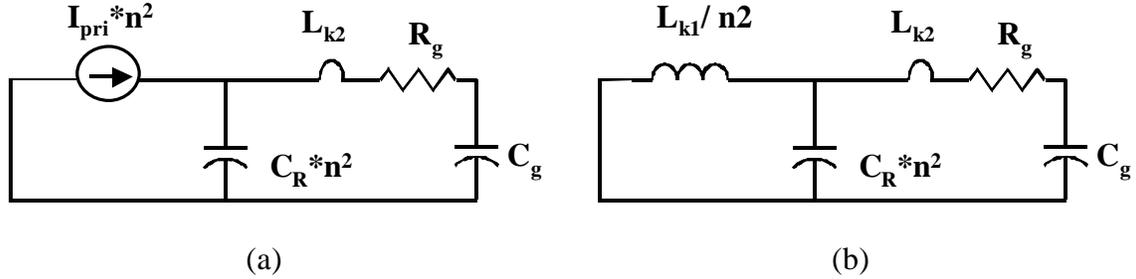


Fig. 2.26. Equivalent circuit diagram for the self-driven circuit.

The circuits shown in Fig. 2.26 were simulated using the SABER simulation tool to analyze the turn-on and turn-off loss savings. The simulation was set up with the following conditions: the operating frequency is 1MHz, the input voltage is 48V, the power transformer has a 10:1 turns ratio, the gate driver transformer has a 5:1 turns ratio, the load current is 80A, the secondary side leakage inductor is 10 nH, the primary devices are Hitachi's HAT2175 that have a  $C_{oss}$  of 185pF, an additional 1nF resonant capacitor is placed in parallel with the drain to source of the bottom primary switch and there are 5 Hitachi HAT2165 synchronous rectifiers in parallel. An individual HAT2165 has a  $C_{gs}$  of 5.208 nF at a  $V_{gs}$  of 9.6V and an internal gate resistor of 0.5  $\Omega$ . The internal gate resistor is in series with the gate drive resistor. Because 5 HAT2165 devices are used in parallel, the equivalent resistance is 0.1  $\Omega$  and the total gate capacitance is equal to 26.04nF. Fig. 2.27 shows the turn-on losses of the self-driven circuit vs. a conventional external voltage driver for one leg of the current doubler. The simulation was run while varying the gate drive resistor  $R_g$  and the gate drive transformer leakage inductor.

Seen in Fig. 2.27, the self-driven circuit can save gate driver turn-on losses. For example, an additional 0.25  $\Omega$  gate drive resistor with 1nH of leakage inductance can save 45% of the turn on gate driver losses.

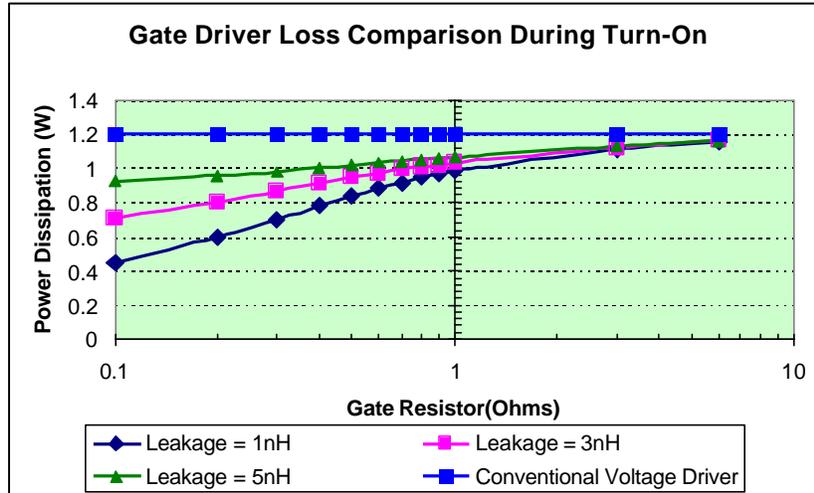


Fig. 2.27. Turn on loss comparison of the self-driven circuit vs. a conventional external voltage driver, both operating at 1MHz for 5 HAT2165 in parallel.

Fig. 2.28 shows the turn-off losses of the self-driven circuit vs. a conventional external voltage driver. The simulation was run while varying the gate drive resistor  $R_g$  and the gate drive transformer leakage inductor.

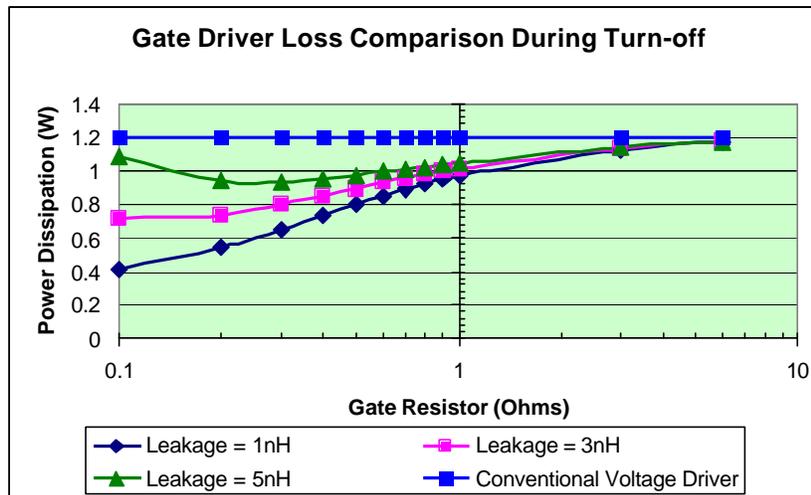


Fig. 2.28. Turn-off loss comparison of the self-driven circuit vs. a conventional external voltage driver, both operating at 1MHz for 5 HAT2165 in parallel.

Seen in Fig. 2.28, the self-driven circuit can save gate driver turn-off losses. For example an additional  $0.25 \Omega$  gate drive resistor with 1nH of leakage inductance can save 50% of the turn-off gate driver losses. During turn-off, the equivalent gate drive circuit is a resonant circuit. The gate driver losses below a  $0.2 \Omega$  gate resistor actually begin to

increase, but remain below the conventional external driven circuit. For example in the 5nH leakage inductor case the losses at 0.1  $\Omega$  are about 1.1W while 0.95W at 0.25  $\Omega$ .

During a turn-off transition, the leakage inductor  $L_{k2}$ , shown in Fig. 2.25, is charged by the gate capacitance of the synchronous rectifiers. Once the synchronous rectifier gate to source voltage drops below the forward voltage of the source-to-gate diode, the leakage inductor will charge  $C_S$  beyond the secondary side of the gate drive transformer voltage. If the leakage inductor charges  $C_S$  beyond the secondary side of the gate drive transformer voltage, the difference can possibly reach the threshold voltage of the synchronous rectifiers. If the synchronous rectifier turns back on, there will be a simultaneous conduction of both synchronous rectifiers and a short circuit will occur. An example of this process is shown in Fig. 2.29.

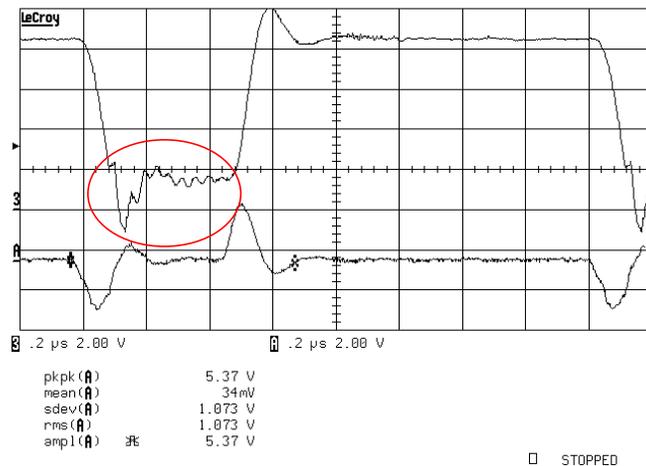


Fig. 2.29. Test waveforms of a self-driven complementary-controlled full bridge operating at 650 kHz. (top waveform)  $V_{GS}$  of a synchronous rectifier, (bottom waveform) voltage across a 0.5  $\Omega$  gate drive resistor.

Depicted in Fig. 2.29 are test waveforms of a self-driven complementary-controlled full bridge operating at 650 kHz. The synchronous rectifiers are 5 HAT2165 from Hitachi placed in parallel. Circled in the top waveform of Fig. 2.29 is the turn-off period of the synchronous rectifier. The gate voltage drops below negative 0.7 V and then is brought back to near positive 2-volts. For example, if the threshold voltage of the synchronous rectifier is less than 2-volts, then the switch will turn back on.

Another possibility for the  $V_{C_2}$  to be greater than the secondary side gate drive transformer voltage is during a transient. In a steady state operation,  $V_{C_S}$  will follow the reflected voltage of  $C_P$ . The voltage of  $C_P$  given in (2.6) is duty cycle and input voltage controlled. During turn off, the voltage of  $C_P$  controls the secondary side gate drive transformer voltage. In a steady state operation, the secondary side gate drive voltage equals  $V_{C_S}$  that turns off the SR's. If  $V_{C_S}$  is less than the secondary side gate drive transformer voltage,  $C_S$  is charged through a forward biased diode to equal the secondary side gate drive voltage. During a transient, if the load current is increased or the input voltage is decreased, the duty cycle will enlarge. Then the secondary side gate drive transformer voltage will decrease to less than  $V_{C_S}$ . Because there is not a low resistance path for  $C_S$  to discharge during a transient, a net positive voltage will be left at the gate of the synchronous rectifiers. Therefore, if the voltage difference is greater than the threshold voltage of the synchronous rectifier, then the synchronous rectifier will not turn off causing a secondary side shoot through current.

To make  $C_S$  rapidly track  $C_P$ , a low resistance path is needed. Therefore, the diode in Fig. 2.25 is replaced by a tiny MOSFET, shown in Fig. 2.30. An additional winding is coupled to the gate drive transformer to control the tiny MOSFET.

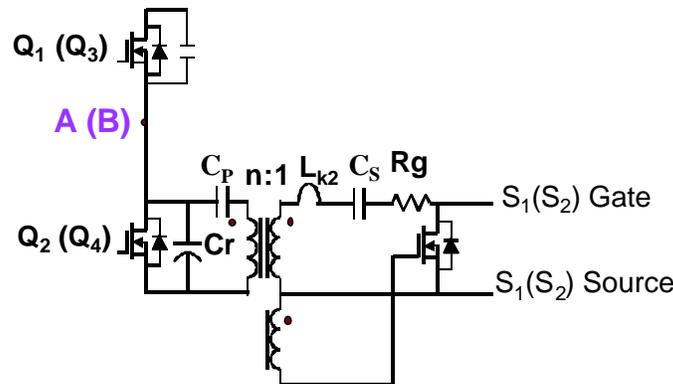


Fig. 2.30. Self-driven circuit with the additional tiny MOSFET.

With the addition of the tiny MOSFET,  $C_S$  will track  $C_P$ , but now a path is formed for the leakage inductor  $L_{k2}$  to freely oscillate, as shown in Fig. 2.31. Circled in Fig. 2.31

is the leakage inductance oscillation, as the tiny MOSFET turns on. The potential problem with having a fast turn on of the tiny MOSFET is that the resonant peak of the oscillations could cross the threshold of the synchronous rectifier. If the threshold is reached, there will be a secondary side shoot through current. Therefore, the timing of the tiny MOSFET is critical.

If the tiny MOSFET does not turn off before the synchronous rectifiers are to turn on, then the synchronous rectifier turn on signal is shorted to ground through the tiny MOSFET. In this case, because the turn on of the synchronous rectifiers is under a virtual ZCS situation, the synchronous rectifier turn on speed is not too critical.

If the tiny MOSFET turns on before the ZVS action of the primary switches completely discharges the gate of the synchronous rectifiers, the gate drive loss savings will be reduced. Likewise, under a non-ZVS condition having the tiny MOSFET turn on quickly will reduce the potential for a secondary side shoot through current due to a slow SR turn off speed.

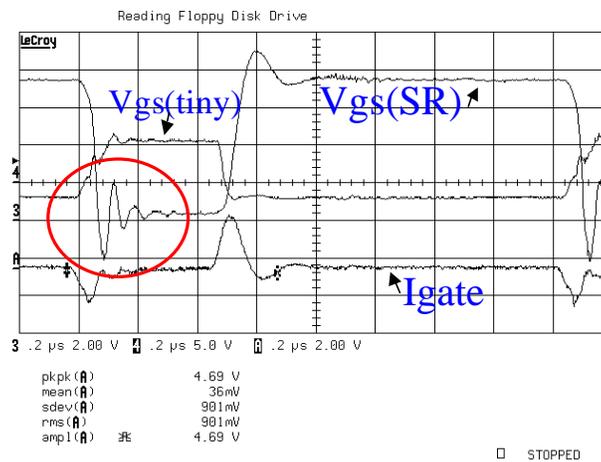


Fig. 2.31. Test waveforms of a self-driven complementary-controlled full bridge operating at 650 kHz. (top waveform)  $V_{GS}$  of a synchronous rectifier, (middle waveform)  $V_{GS}$  of the tiny MOSFET, (bottom waveform) voltage across a  $0.5 \Omega$  gate drive resistor.

To completely recover all of the possible gate drive losses, a delay must be added in the tiny MOSFET gate drive path. To allow for the tiny MOSFET to turn off quickly, there must be a low resistance discharge path. Therefore, the following gate drive circuit, seen in Fig. 2.32, was designed to take advantage of these properties. A 200-Ohm

resistor is added in the gate drive path to slow down the turn on and a Schottky diode is placed in parallel to achieve a fast turn-off. The tiny MOSFETs used are Si3900 from Siliconix Inc.

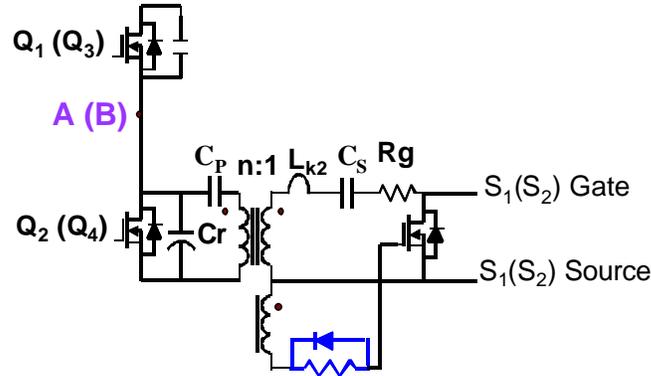


Fig. 2.32. Self-driven gate drive circuit with a delay in the tiny MOSFET gate drive path.

Seen in Fig. 2.33 are the operational waveforms of the self-driven circuit shown in Fig.2.32. Circled are the gate drive turn-on and turn-off of the tiny MOSFET and the turn-off time of the synchronous rectifiers. The slope of the turn-on of the tiny MOSFET is very slow and the turn-off is very rapid. There are nearly no oscillations during the turn-off time of the synchronous rectifiers. Also shown in Fig. 2.33 is the gate drive current. The peak and shape during the turn-off of the synchronous rectifiers is nearly equivalent to the turn-on current. Having the peak and shape equivalent shows that there is no gate drive energy dissipated in the tiny MOSFET.

Using the gate driver shown in Fig. 2.32 has the advantages of recovering the most possible gate drive energy. However, the addition of a slow tiny MOSFET can be non-advantageous under light load or non-ZVS conditions. The leakage inductor in the main power transformer dictates the timing of the tiny MOSFET.

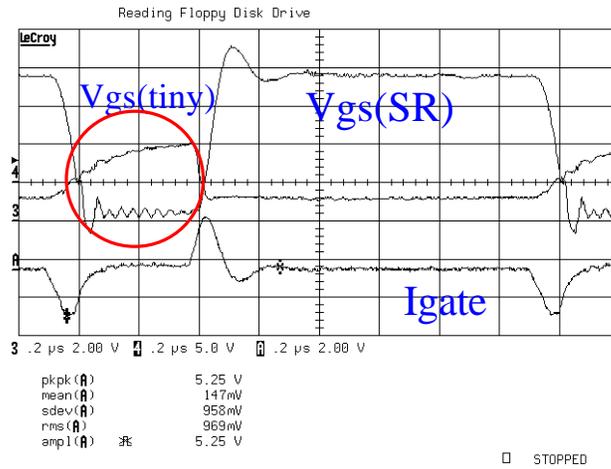


Fig. 2.33. Operational waveforms of the self-driven circuit shown in Fig. 2.32.

The current in the main power transformer passes from a positive to a negative direction, or vice-versa, to begin each power delivery period, shown in Fig. 2.18. The slope of this transition is given by (2.11).

$$slope = \frac{V_{IN}}{L_K * n^2} \quad (2.11)$$

Where  $L_K$  is the leakage inductance referenced to the secondary side of the transformer and  $n$  is the turns ratio.

If the synchronous rectifier turns off after the main power transformer current changes direction, the current delivery path on the secondary side will be a short circuit. Therefore, there will be a shoot through current and the losses will rapidly increase.

If the synchronous rectifier turns off before the transformer current has changed directions, then the body diode will conduct. The time in which the transformer current changes direction is as follows:

$$\Delta t = \frac{\frac{I_o}{2} + \Delta i_o - \frac{V_o}{(L_{LK} + L_{F1})} (1-D) * T_s}{\frac{V_{IN}}{nL_{LK}}} \quad (2.12)$$

where  $\Delta i_O$  is the change in output filter current,  $L_{LK}$  is the leakage inductor seen on the secondary side,  $L_{LF1}$  is one filter inductor,  $D$  is the duty cycle, and  $T_s$  is the switching period.

Fig. 2.34 compares the operational waveforms between phase-shifted and complementary-controlled full bridge circuits. Because the self-driven gate drive signals respond directly to the points A and B on the primary side, the self-driven complementary-controlled full bridge has a shorter body diode conduction time than the phase-shifted full bridge. The body diode conduction loss comparison can be seen in Fig. 2.35, under a load condition of 1.2V/80A at 1MHz switching frequency with 10 nH leakage inductor and a 20 ns delay time. The loss saving is 2.56 W, which is a 34% reduction, and corresponds to a 2.5% total efficiency improvement.

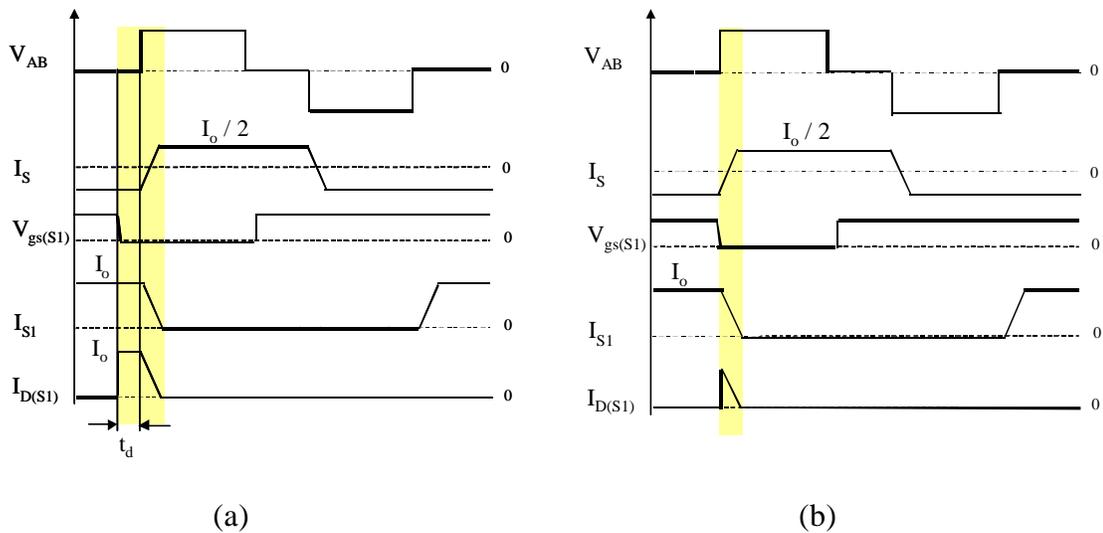


Fig. 2.34. Comparison between the body diode conduction times (a) Phase-shift full bridge, (b) Self-Driven complementary-controlled full bridge. Highlighted are the body diode conduction times.

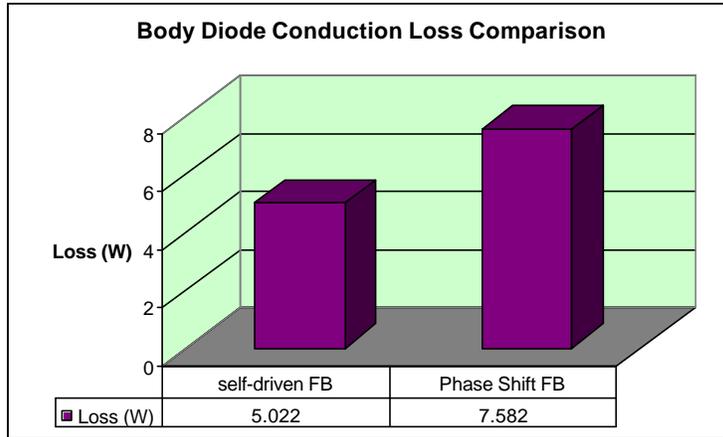


Fig. 2.35. Body diode conduction loss comparison.

Fig. 2.36 compares the efficiency of the self-driven complementary-controlled full-bridge with a fast tiny MOSFET and delayed tiny MOSFET depicted in Fig. 2.32. As can be seen the efficiency under the light load conditions (pre-ZVS) the converter with a fast tiny MOSFET has improved efficiency. After ZVS is achieved, both converters have nearly identical efficiencies.

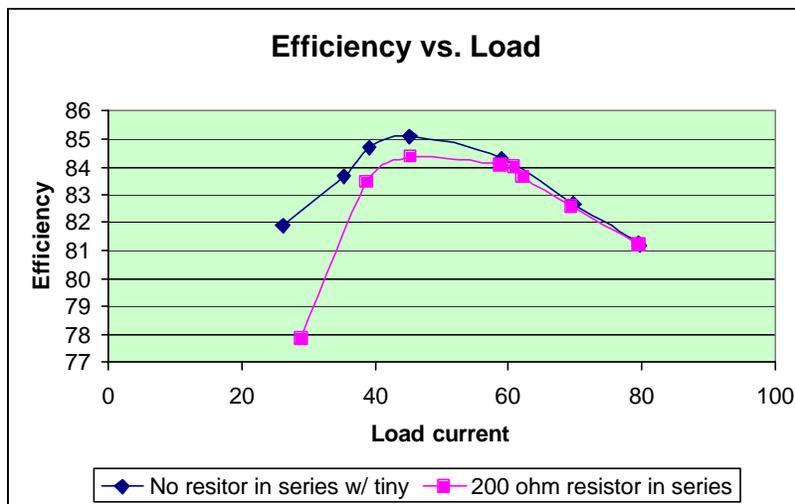


Fig. 2.36. Efficiency comparison of the self-driven circuit with a fast tiny MOSFET and a delayed tiny MOSFET operating 650 kHz with an output of 1.2V and an input of 48V.

Therefore, to improve the light load efficiency a faster turn-off speed for the synchronous rectifiers is needed.

## 2.7 Summary

This chapter presents and analyzes a self-driven gate drive concept used on a complementary-controlled full bridge circuit. The self-driven circuit shows a great improvement over the conventional external driven circuit. The simplicity, loss savings and small size of the self-driven topology lends itself to compact high-frequency isolated DC/DC conversion. With a 0.25-Ohm gate drive resistor and an operating frequency of 1MHz, while driving 5 HAT2165 synchronous rectifiers, on each leg of the current doubler, there is almost 50% gate driver loss savings compared to a conventional external driver. This is a 2.5 W reduction. When comparing the phase-shift to the self-driven complementary-controlled full bridge, the body diode loss savings are about 34%, a 2.56 W savings. Therefore, using the self-driven complementary-controlled full bridge for a 48V to 1.2V/80A, operating at 1MHz, there is a 5.1% efficiency improvement over the phase-shifted full bridge operating under the same conditions.

The presented self-driven topology has the following beneficial attributes:

- The self-driven topology holds the synchronous rectifiers to the full gate drive voltage during dead times on the power transformer
- The input gate capacitance of the synchronous rectifiers operates a lossless snubber capacitor for the primary side switches
- 50% gate drive savings compared to an external driver
- Minimized synchronous rectifier body diode conduction time
- Very clean synchronous rectifier gate drive signals due to tight coupling of the gate drive windings.
- The gate drive windings are not coupled to the power windings. Therefore, the power windings can be optimally wound, which minimizes the leakage inductor and the secondary side winding resistance by utilizing all layers of the PCB for power windings. Also there are discrete

gate drive transformers that do not allow the power winding leakage current to couple into the gate drive windings.



## **Chapter 3**

# **Performance Improvements of Integrating Multiple Transformers**

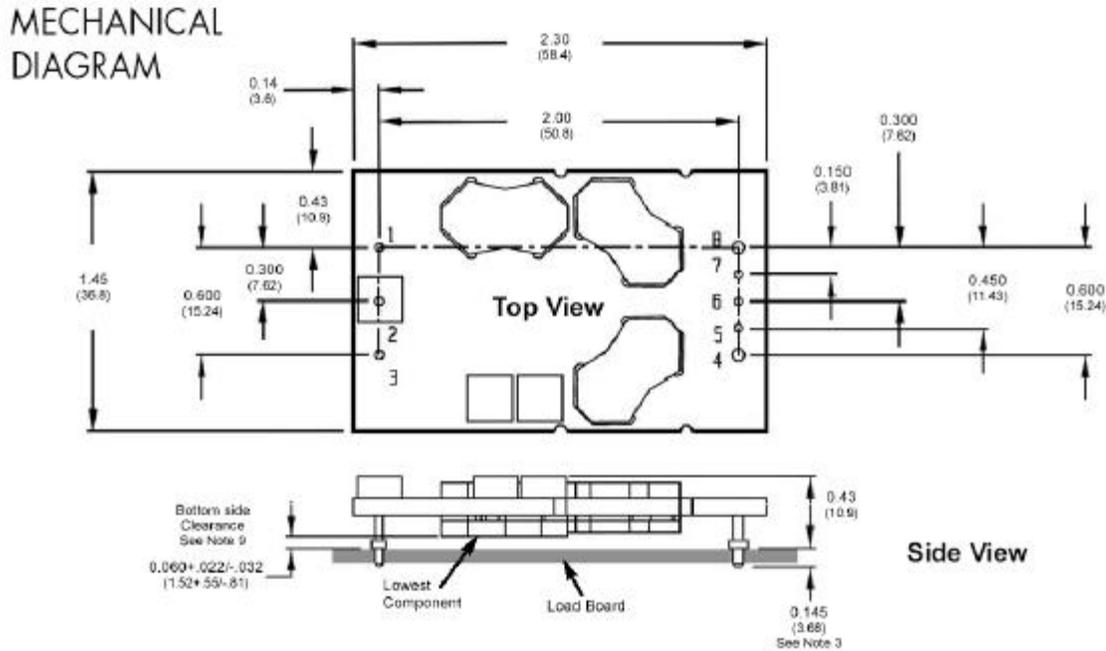
The previous chapter illustrates a novel technique for adapting a self-driven scheme to the complementary-controlled full bridge topology. However, with the addition of two gate driver transformers, the implementation of the self-driven circuit does not lead to an overall compact design. This chapter presents an integrated transformer concept that can be used to reduce circuit complexity, the overall size and cost of the power stage. In addition, this concept increases the efficiency of the circuit.

### **3.1 Motivation**

In Chapter 2, the self-driven scheme used to drive the synchronous rectifiers was adapted to the complementary-controlled full bridge. The implementation of the self-driven circuit gave an additional two transformers that increased the circuit's magnetic core number to five, of which there are three transformers and two output filter inductors. Normally, a designer uses a self-driven circuit not only for the loss saving abilities but also for the size reduction. The addition of two gate drive transformers can reduce the attractive properties of the self-driven circuit discussed in Chapter 2.

In a compact DC/DC converter, size is always a driving factor and the thermal dissipation is always the limiting factor. A modular DC/DC converter, such as a VRM or brick type converter, has a fixed form factor. For example, in designing a quarter brick all of the electrical requirements must be placed into a standardized 1.45 inches x 2.3

inches x 0.452 inches for the width, length and height, respectfully. The quarter brick also has isolation requirements, temperature limitations, standard input and output pin placements, as well as many other design specifications. The mechanical diagram for a PowerQor Tera PQ60012QTA40 Quarter-Brick from Synqor is shown in Fig. 3.1 and as seen the placement of the transformers and output filter inductors is critical.



(From Synqor's PQ60012QTA40 datasheet [www.synqor.com](http://www.synqor.com))

Fig. 3.1. Mechanical Diagram of a PowerQor Tera PQ60012QTA40 Quarter-Brick from Synqor.

Shown in Fig. 3.1, there are three large magnetic components. Synqor's patented topology shown in Fig. 3.1, is a two-stage design where the first stage acts as a current source and the second stage is a self-driven push-pull topology with two transformers that form a current doubler like rectifier. Synqor's topology uses a form of self-driven synchronous rectification that cross-couples the synchronous rectifier gate drive signals from the two output transformers. In order to use the self-driven scheme outlined in Chapter 2, there must be an additional two gate drive transformers. Increasing the switching frequency can decrease the physical size of the magnetic components, but the losses and ability to carry the large current demand of future brick type converters can limit the overall reduction in size. In a quarter-brick design the magnetic components are

typically of the planar type to meet the modules height restrictions. When placing a planar magnetic component into a printed circuit board (PCB), large holes will be made through the board to fit the magnetic component. These holes will reduce the area in which the feedback signals and power traces can be placed.

One possible arrangement for the main power stage components of the self-driven complementary-controlled full bridge is shown in Fig. 3.2. The placement of the power transformer must be as close to the output as possible to limit the high current secondary copper losses. Other than the sheer size of the magnetic components, the major problem with this design is that the gate drive length from the secondary of the gate drive transformer to the synchronous rectifiers is a long path.

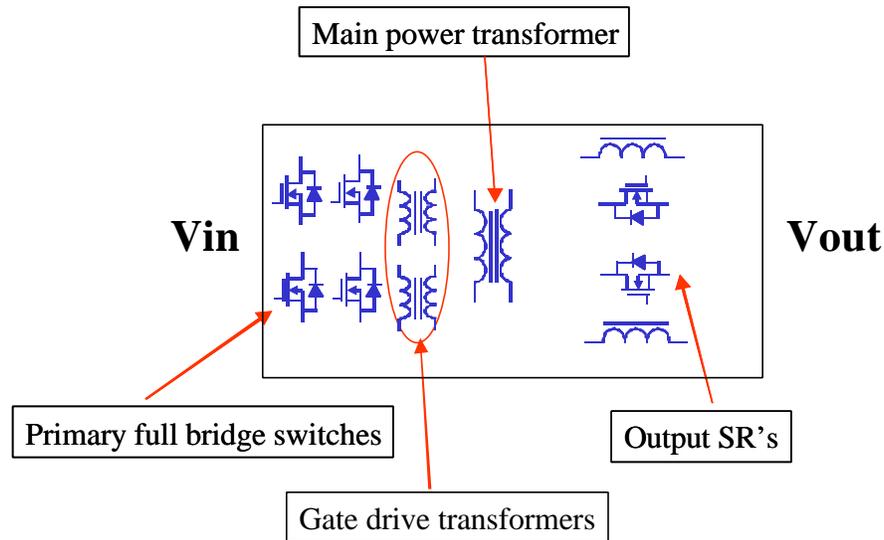


Fig. 3.2. A topological arrangement for the three gate drive transformers and two output inductors.

Having a long secondary side gate drive length will introduce larger parasitic trace inductance and impedance that can slow down the switching speed of the synchronous rectifiers, as well as limit the loss savings. In a compact design it is not feasible to use a three transformer and two output inductor design. Using the current technologies of integrating the output inductors with the main power transformer can eliminate the need for the two output inductor cores, but this does not solve the problem of the long gate drive length. To make this self-driven scheme more attractive in a

compact high frequency DC/DC converter design, the gate drive transformers must push to the secondary side of the power transformer, or be integrated directly with the power windings.

This chapter presents a new concept for magnetic integration by integrating three transformers onto a single magnetic core. Each secondary winding such as the power secondary and the two gate driver secondary windings will directly couple to a discrete set of primary winding. This allows each secondary/primary winding pair to use a properly interleaved winding structure that yields the smallest leakage inductance possible. Presented is the derivation of the integrated transformer, flux path in the integrated transformer and methods to calculate the losses. Two prototypes were constructed in the quarter brick form factor to validate the integrated transformer techniques, one was built with three discrete transformers and the second was built with the integrated transformer. After the first two prototypes validated the concept, a third prototype was designed to take advantage of the benefits of the integrated transformer to increase the converters output current and efficiency.

### **3.2 Derivation of the Integrated Transformer**

The complementary-controlled full bridge, Fig. 3.3, has four major functional states, two where power is delivered, and two that are freewheeling state. The operational waveforms of the power transformer and point A and point B for these four states are shown in Fig. 3.4.

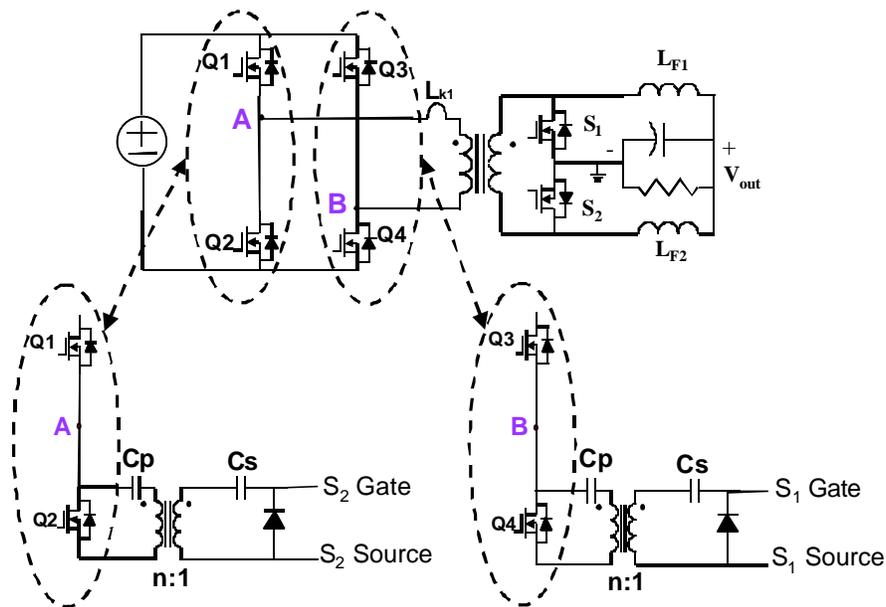


Fig. 3.3. Complementary-Controlled Full Bridge with Self-Driven Synchronous Rectifiers.

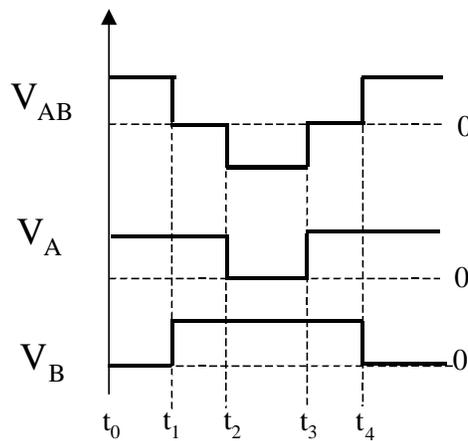


Fig. 3.4. The operational waveforms of the power transformer and point A and point B.

Referring to the waveforms in Fig 3.4 and the equivalent circuit topology in Fig. 3.5 from time  $t_0$  to time  $t_1$ , the switches  $Q_1$  and  $Q_4$  are turned on and the self-driven circuit holds  $S_2$  on and  $S_1$  is off. The voltage at point A is equal to the bus voltage and point B is pulled to ground by  $Q_4$ . The voltage polarities across each transformer are shown in Fig. 3.5a. Fig. 3.5b shows the resulting AC flux directions that occur in the

corresponding transformer. The main power transformer has a net AC flux given by Faraday's Law shown in (3.1) and (3.2). The main power transformer is shaped as an EE/EI core and, therefore, the net flux in the center leg splits evenly between the two outer legs.

$$N \cdot \mathbf{f} = \int \mathbf{v} \cdot d\mathbf{t} \quad (3.1)$$

$$V = n * \frac{d\mathbf{f}}{dt} \quad (3.2)$$

Where V is the bus voltage, n is the number of primary turns and  $\phi$  is the flux.

The two gate drive transformers also have a net AC flux in the cores. The driver transformers A and B are shown as a UU/UI core. Because the voltages applied across the two gate drive windings are opposite in polarity, the net AC flux will propagate in opposite directions. The voltage across drive winding A is equal to the voltage at point A minus the voltage of  $C_{P1}$ . The voltage across drive winding B is equal to the voltage of  $C_{P2}$ .

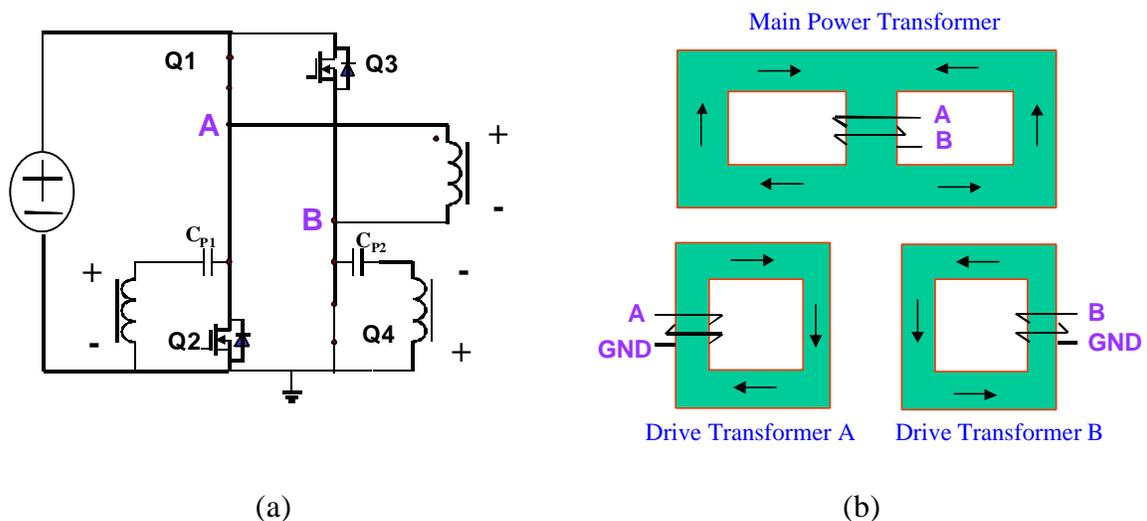


Fig. 3.5. From time  $t_0$  to  $t_1$  (a) First power delivery topological state (b) Resulting flux in the three transformers. The arrows state the direction of the AC flux in the cores.

The time period from  $t_1$  to  $t_2$  is the first of the two freewheeling periods. At the transition from  $t_1$  to  $t_2$ , the switch  $Q_4$  turns-off and the switch  $Q_3$  turns-on, resulting in  $S_1$

turning on. Switches  $Q_1$  and  $S_2$  remain on. The voltage across the main power transformer windings is zero and the voltage across drive transformer B changes direction. The voltage at point A remains at the bus voltage, while point B is pulled up to the bus voltage. The flux in drive transformer A will continue to flow in the same direction. Because the voltage across the drive transformer B changes polarity, the resulting flux will change direction. During the freewheeling states, net flux in the main power transformer is zero. The topological state and resulting transformer fluxes are shown in Fig. 3.6.

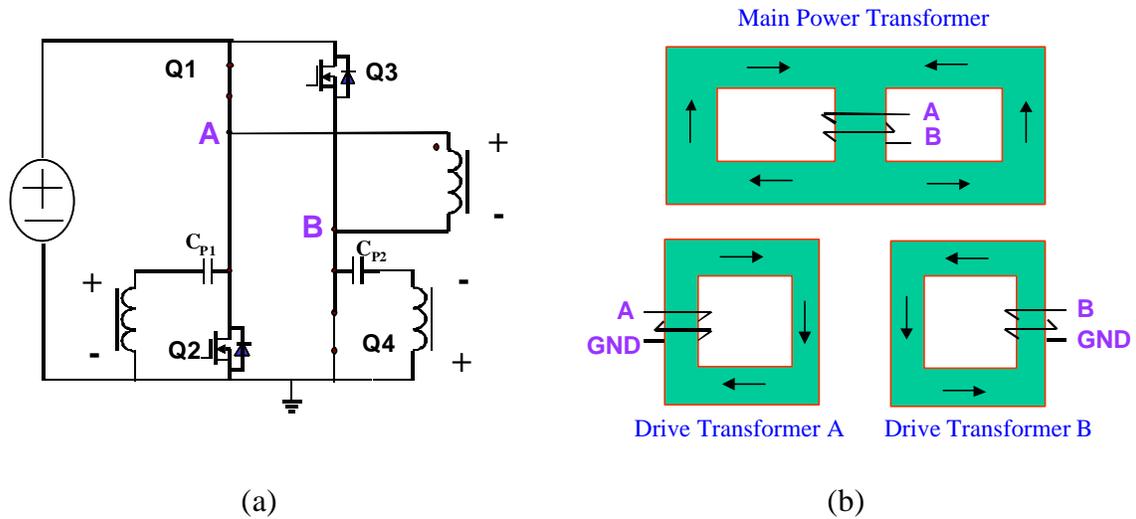


Fig. 3.6. From time  $t_1$  to  $t_2$  (a) Second freewheeling topological state (b) Resulting flux in the three transformers. The arrows state the direction of the AC flux in the cores.

With respect to the polarities on the transformers the time period from  $t_2$  to  $t_3$  has exact opposite operation as from time  $t_1$  to  $t_2$ . Switches  $Q_3$  and  $S_1$  remain on, while  $Q_1$  and  $S_2$  turn-off and  $Q_2$  turns-on. The voltage at point B remains at the bus voltage and point A is pulled to ground. Therefore, there is a net voltage that is applied to the main power transformer that results in a net flux in the core. The flux will propagate in the opposite direction, as from time  $t_1$  to  $t_2$ . The flux in driver transformer B will remain flowing in the same direction as the previous topological stage, while the flux in driver transformer A will change direction resulting in  $S_2$  turning off. The topological state and resulting transformer fluxes are shown in Fig. 3.7.



From examination of the flux directions in all three transformers over the entire switching period, it is apparent that the flux direction in drive transformer A and B map directly to the left and right side of the main power transformer, respectfully. During the two power delivery periods there is a net flux produced in the center leg of the EE/EI core. As pictured in Fig. 3.5 and 3.7, the flux direction in the two driver transformers both flow in the same direction as the flux in the power transformer. During the two freewheeling periods there is zero net flux in the center leg of power transformer. The flux in the two gate drive transformers corresponding to the center leg in the power transformer flow in opposite directions. The flux produced in the gate drive transformers during the freewheeling period is as follows:

$$\frac{d\mathbf{f}_A}{dt} = \frac{V_{IN} - V_{CP1}}{n} \quad , \quad \frac{d\mathbf{f}_B}{dt} = \frac{V_{IN} - V_{CP2}}{n} \quad . \quad (3.3)$$

Where  $\phi_A$  and  $\phi_B$  are the flux produced in the individual gate drive transformer, and  $n$  is the turns ratio of the transformer. Because each transformer has the same turns ratio, and the voltage  $V_{CP1}$  equal  $V_{CP2}$ , the flux produced by each transformer are identical. By, winding each gate drive transformer in the opposite direction and mapping the gate drive transformers to the left and right sides of the power transformer during the freewheeling period, the net flux produced in the center leg is zero.

Following the above logic the integrated transformer can be derived. Starting from the three discrete transformers the derivation of the integrated transformer is shown in Fig. 3.9. An EE/EI core is used, and the main power winding, drive winding A and drive winding B, are place on the center leg, on the left and right leg, respectfully.

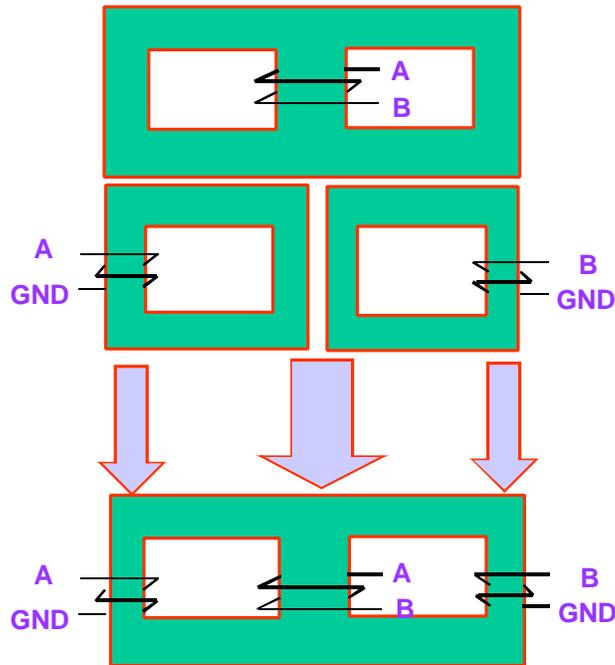


Fig. 3.9. Integrated transformer windings.

From the integrated transformer winding structure shown in Fig. 3.9, the secondary power and gate drive windings can be coupled directly to the corresponding primary winding as shown in Fig. 3.10. Each secondary winding having a discrete primary winding allows for a fully interleaved winding structure that results in the lowest obtainable leakage inductance for each set of windings.

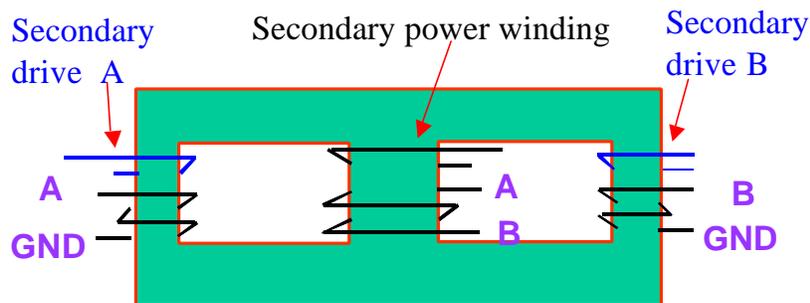


Fig. 3.10. Integrated transformer winding structure.

At first glance, one would believe that this is not feasible or that the core would need to be gapped. The following analysis will show that this is possible and that the core does not need to be gapped.

### 3.3 Analysis of the Integrated Transformer Structure

Beginning from a single winding structure, Fig. 3.11, a common inductor can be realized. The AC flux depends only on the AC voltage applied, given by Faraday's Law. The inductive winding will be known as an active winding because voltage is applied directly across the winding. An example of a two active winding structure would be a coupled inductor.

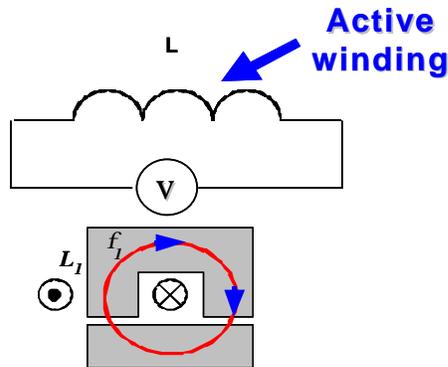


Fig. 3.11. Basic inductor.

A two winding core structure can have an active winding and also a passive winding, i.e. a transformer, Fig. 3.12. A voltage is applied to the primary, or active winding and the secondary, or passive winding, will respond to the action resulting from the primary. According to Faraday's Law and Lenz's Law, if the flux induced on the secondary winding is known, then the resulting secondary voltage is given by (3.4).

$$V_{SEC} = \frac{N_s}{N_p} * \frac{d\mathbf{f}}{dt} \quad (3.4)$$

Where  $V_{sec}$  is the voltage produced across the secondary winding,  $N_s$  is the number of secondary windings,  $N_p$  is the number of primary windings,  $d\phi$  is the change in flux and  $dt$  is the change in time.

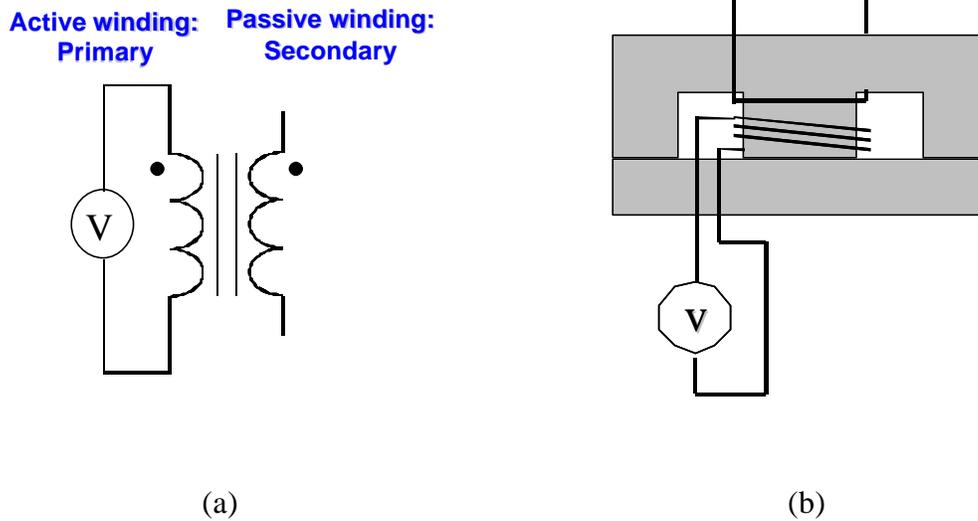


Fig. 3.12. Basic Transformer (a) circuit implementation, (b) physical implementation.

A transformer can also have multiple primary active windings seen in Fig. 3.13 as long as (3.5) is satisfied.

$$\frac{V_1}{N_1} = \frac{V_2}{N_2} \quad (3.5)$$

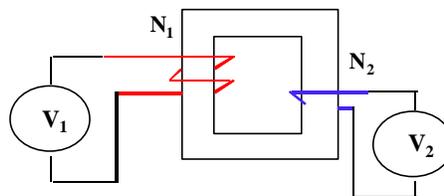


Fig. 3.13. Multiple primary winding structure.

If in (3.5) the left side of the equation is greater than the right or vice-versa, the current from the greater value will increase without bound to infinity. For example, assume  $V_2$  is a short circuit and  $V_1$  is some arbitrary voltage. The flux imposed on the

core from  $V_1$  will be some value  $\phi_1$  given by Faraday's Law, where  $\phi_1$  does not equal 0, and the flux imposed by  $V_2$  will be  $\phi_2$ , where  $\phi_2$  equals 0. Then the flux  $\phi_1$  will need a path to flow, and because  $V_2$  is a short circuit there is no flux through the windings of  $V_2$ . This means that all of  $\phi_1$  will travel through the air. The reluctance of air is near infinity compared to the reluctance of the core. Therefore, the mmf across  $R_{air}$ , seen in Fig. 3.14, will increase to infinity because  $R_{air}$  is equal to infinity. Consequently, the mmf across the windings will go to infinity.

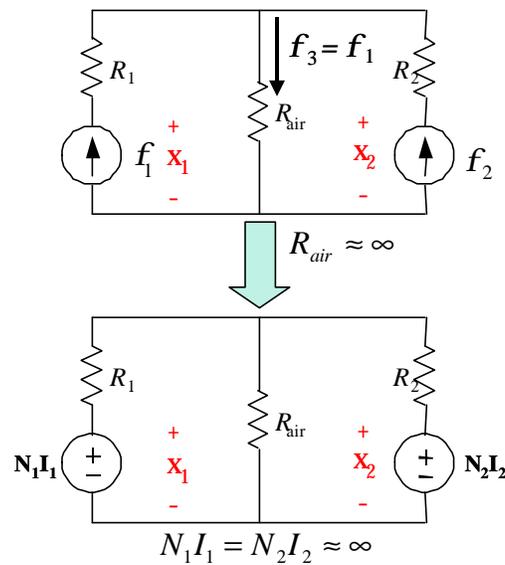


Fig. 3.14. Reluctance models.

It follows from the above analysis that if (3.5) is satisfied, then placing two primary active windings on a single magnetic core is acceptable. By applying the above magnetic concept to the circuit topology given in Fig. 3.3, the following self-driven complementary-controlled full bridge using the integrated transformer can be realized, shown in Fig. 3.15.

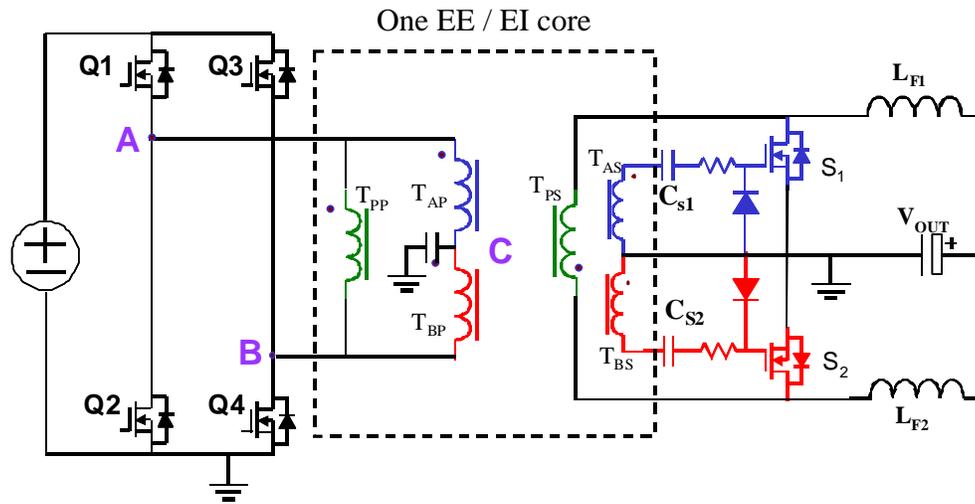


Fig. 3.15. Schematic of the self-driven complementary-controlled full bridge with the integrated transformer.

On the primary side of the circuit shown in Fig. 3.15, there is only one gate drive capacitor that replaces both  $C_p$ 's in Fig. 3.3. Because the voltages across both  $C_p$  capacitors are identical, the second gate drive capacitor is redundant. The winding structure for the self-driven complementary-controlled full bridge with the integrated transformer is shown in Fig. 3.16 and the primary side transformer waveforms are shown in Fig. 3.17. In order to have (3.5) balance, all of the primary side winding for each the power and the gate drive windings must have the same number of turns.

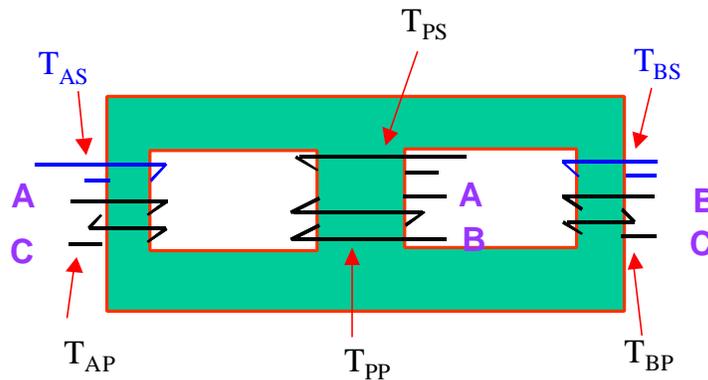


Fig. 3.16. Winding structure for the self-driven complementary-controlled full bridge with the integrated transformer.

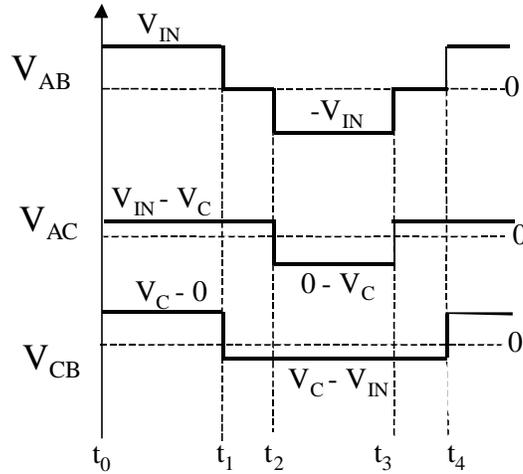


Fig. 3.17. Primary side transformer waveforms for the self-driven complementary-controlled full bridge with the integrated transformer.

Referring to Fig. 3.15, 3.16 and 3.17, the secondary power winding  $T_{SP}$  is tightly coupled to the power primary winding  $T_{PP}$ . The two gate drive windings also have a corresponding tightly coupled secondary winding,  $T_{AP}$  to  $T_{AS}$  and  $T_{BP}$  to  $T_{BS}$ . Having a primary winding tightly coupled to the corresponding secondary winding using proper interleaving techniques will yield the least amount of leakage inductance possible in each leg. A discussion of the operation of this circuit follows.

From  $t_0$  to  $t_1$  power is transferred to the secondary side. The winding from A to B has a positive voltage applied and a net flux circulates in the center leg of the transformer. Winding  $T_{AS}$  will apply a positive voltage to the gate of  $S_1$  and winding  $T_{BS}$  will turn off  $S_2$ .

From  $t_1$  to  $t_2$ , power will no longer flow to the secondary side. Switches  $Q_1$  and  $Q_3$  will both be on, allowing current to freewheel. Since,  $Q_1$  and  $Q_3$  are both conducting, the bus voltage will be applied to windings  $T_{AP}$  and  $T_{BP}$ . Winding  $T_{AP}$  will cause clockwise flux path, depending on winding direction, in the left hand side of the transformer and winding  $T_{BP}$  will cause a clockwise flux in the right hand side of the transformer of the same magnitude as  $T_{AP}$ . Thus, the center leg will have the flux created by  $T_{AP}$  minus the flux created by  $T_{BP}$ , which will cause the sum of the flux to be equal to zero.

From  $t_2$  to  $t_3$ , switches  $Q_2$  and  $Q_3$  are conducting. The bus voltage will be applied to the  $T_{PP}$  in the opposite direction as  $t_0$  to  $t_1$ . The winding voltage across  $T_{AP}$  will be the negative capacitor voltage causing the polarity of  $T_{AS}$  to change direction. This results in turning S1 off. Winding  $T_{BP}$  will remain having a positive voltage across the windings, holding S2 on.

The operation time from  $t_3$  to  $t_4$  is exactly the same as from time  $t_1$  to  $t_2$ . From time  $t_4$  on another cycle occurs.

### 3.4 Loss Calculations in the Integrated Transformer

Beginning from a basic discrete transformer only showing the primary side windings, the following structure can be derived, shown in Fig. 3.18a. If a voltage is applied from point A to B, then following (3.2) a flux,  $\Phi_1$ , is produced in the center leg. The sum of the flux produced in the two outer legs equals the flux produced from the center leg. From the analog of KCL

$$\Phi_1 = \Phi_2 + \Phi_3 \quad (3.6)$$

Most common EE/EI cores have the two outer legs cross-sectional area equal to  $\frac{1}{2}$  of the center legs, as depicted in Fig. 3.18b. Then, if the core is symmetrical  $\Phi_2$  will equal  $\Phi_3$  shown in Fig. 3.18a. To calculate the losses in the transformer core in Fig. 3.18a first calculate the volt-seconds applied to the core, shown in (3.7). Then the volt-seconds are used to calculate the AC flux density of the center leg, shown in (3.8). Using the material properties given in the datasheet of the core along with the core volume, the losses can be easily calculated.

$$I_1 = \int_{t_1}^{t_2} v_1(t) dt \quad (3.7)$$

$$B_{\max} = \frac{I_1}{2n_1 A_c} \quad (3.8)$$

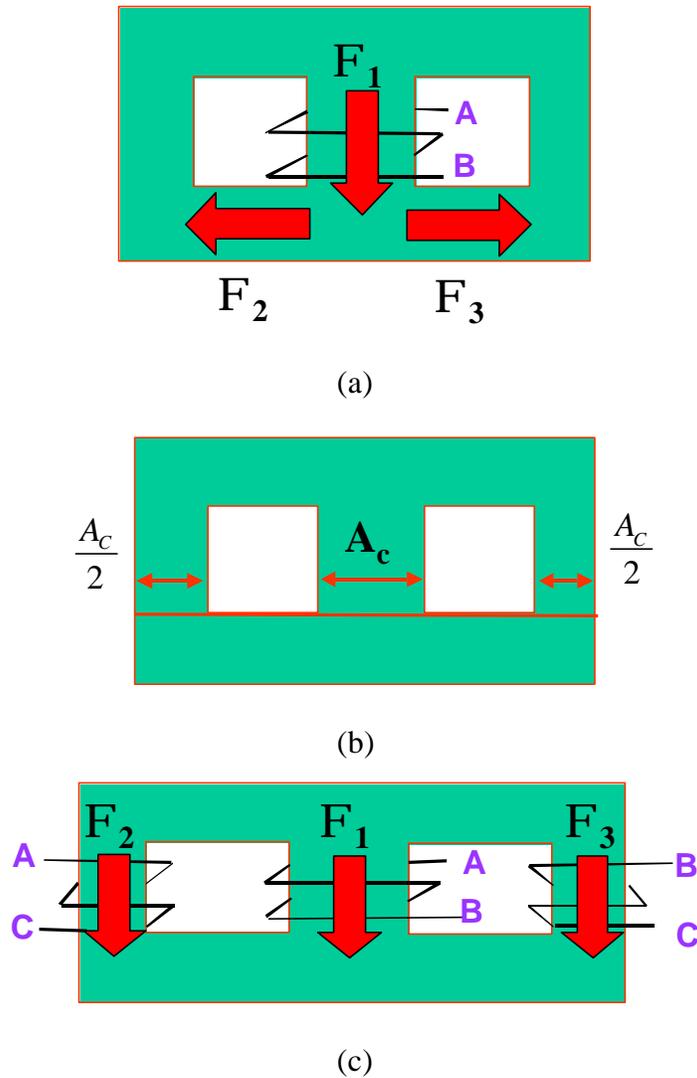


Fig. 3.18. (a) Diagram of a simple transformer (b) Sketch of the cross-sectional area of an EE/EI core, (c) Flux diagram of integrated transformer showing primary side windings.

Fig. 3.18c shows the addition of the two gate drive windings to the outer legs of the transformer. By applying (3.2) to Fig. 3.18c the following are derived

$$V_A - V_B = N_1 \frac{d\Phi_1}{dt} \quad , \quad V_A - V_C = N_2 \frac{d\Phi_2}{dt} \quad , \quad V_C - V_B = N_3 \frac{d\Phi_3}{dt} \quad (3.9)$$

where  $N_1$ ,  $N_2$  and  $N_3$  equal the amount of turns. If  $N_1$ ,  $N_2$  and  $N_3$  are equal then (3.6) can be applied to (3.9) and the flux in the core will balance.

During an entire switching cycle for the self-driven complementary-controlled full bridge there are four major topological changes; two are power delivery periods and

two are freewheeling periods. From a flux point of view, the two power delivery periods are identical with the exception that the flux travels in the opposite direction. For the two freewheeling periods the flux directions are identical. Applying Faraday's Law during one of the two power delivery periods yields the following:

$$\frac{d\Phi_1}{dt} = \frac{V_{in}}{n}, \quad \frac{d\Phi_2}{dt} = \frac{V_{in} - V_C}{n}, \quad \text{and} \quad \frac{d\Phi_3}{dt} = \frac{V_C - 0}{n}. \quad (3.10)$$

And applying Faraday's Law during the freewheeling periods yields the following:

$$\frac{d\Phi_1}{dt} = 0, \quad \frac{d\Phi_2}{dt} = \frac{V_{in} - V_C}{n}, \quad \text{and} \quad \frac{d\Phi_3}{dt} = \frac{V_C - V_{in}}{n}. \quad (3.11)$$

Thus, during the freewheeling periods  $\Phi_2 = -\Phi_3$  and during the power delivery period  $\Phi_1 = \Phi_2 + \Phi_3$ .

The flux in each leg of the transformer must be known to calculate the transformer core losses. The flux produced in the outer legs  $\Phi_2$  and  $\Phi_3$  are a function of the input voltage and the capacitor voltage  $V_{CP}$  (derived in Chapter 2). Because  $V_{CP}$  is a function of the duty cycle given in (3.12), there are two extremes that can occur.  $V_{CP}$  can equal zero or  $V_{IN}$ . At start up  $V_{CP}$  is zero, and then from (3.11), all of the flux will flow through one leg in the transformer. But when a soft-start function is added to the circuit, the value of initial value of  $V_{IN}$  is limited. Therefore, the flux produced from  $V_{CP}$  equal to zero is minimal. Because the circuit topology is a complementary-controlled full bridge, the duty cycle is limited to 50%. Thus,  $V_{CP}$  can never equal zero in any other situation.

$$V_{CP} = V_{IN} * (1 - D) \quad (3.12)$$

In a practical design, a zero duty cycle only occurs when the output voltage is much higher than the reference voltage in the feedback loop such that the compensator cannot control the output. The zero duty cycle will only occur in transient situations. Therefore,  $V_{CP}$  will approach  $V_{IN}$  and, from (3.11), the flux in the core will decrease towards zero.

To calculate the steady state losses of the integrated transformer the volt-seconds of each primary winding must be known. From the volt-seconds, the losses can be calculated in a piece-wise manner. The flux densities for each winding in a symmetrically designed EE/EI core shown in Fig. 3.18b, with the windings shown in Fig. 3.18c, are as follows:

$$\text{Center Leg: } B_{Center} = \frac{1}{n \cdot A_c} \cdot \int_{t_1}^{t_2} V_{AB}(t) \cdot dt \quad (3.13)$$

$$\text{Outside Legs: } B_{Outside} = \frac{2}{n \cdot A_c} \cdot \int_{t_1}^{t_2} V_{AC}(t) \cdot dt \quad (3.14)$$

$$B_{Outside} = \frac{2}{n \cdot A_c} \cdot \int_{t_1}^{t_2} V_{CB}(t) \cdot dt \quad (3.15)$$

The two outer legs flux densities given by (3.14) and (3.15) are different than one half of the flux density of the center leg (3.13). The transformer is divided into three parts shown to calculate the total core losses, in Fig. 3.19. First, the flux densities of the individual parts are calculated by (3.13), (3.14) and (3.15). Then using the material properties of the core given by the manufacturer, along with the volumes of the individual parts as pictured in Fig. 3.19, the sectional losses can be calculated. The total losses are the sum of the losses of the individual parts.

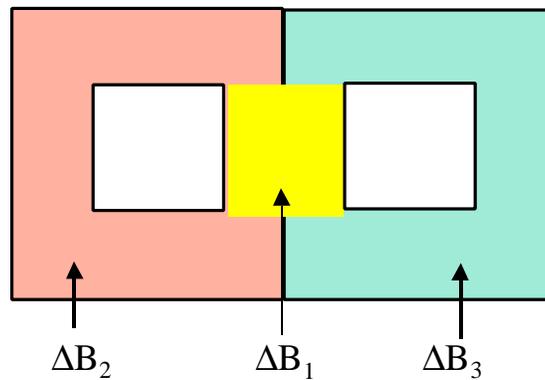


Fig. 3.19. Flux distribution areas.

Maxwell 2D, by Ansoft is used to simulate the integrated transformer to illustrate the flux distribution and core losses shown in Fig. 3.20. The model is setup using a linear ferrite core with the permeability set to 4000. The conductors are a single turn. Each conductor is assigned the respective waveforms derived from the self-driven complementary-controlled full bridge. It is evident from Fig. 3.20 that symmetry exists about the central axis and that flux density is lowest in the center leg of the transformer, verifying the previous theoretical calculations. The flux crowding around the corners of the windows are as expected.

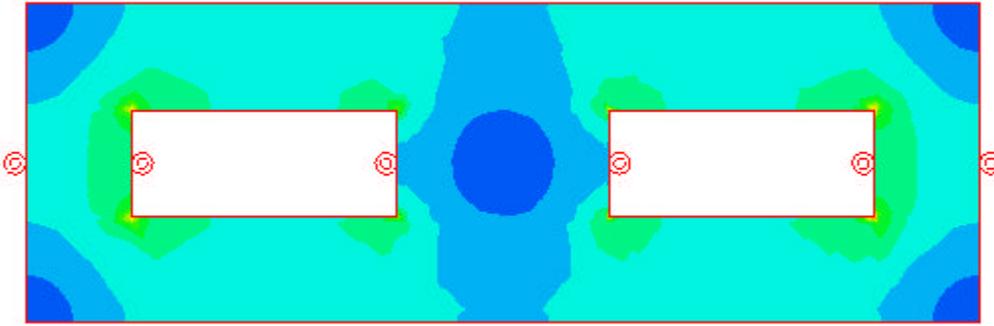


Fig. 3.20. Maxwell 2D simulation of the integrated transformer.

The change in flux  $\Delta B_1$  for the self-driven complementary-controlled full bridge is as follows:

$$\Delta B_1 = \frac{V_{IN} \cdot D}{n \cdot A_c \cdot f_s} \quad . \quad (3.16)$$

The change in flux  $\Delta B_2$  for the self-driven complementary-controlled full bridge is as follows:

$$\Delta B_2 = \frac{(V_{IN} - V_c) \cdot (1 - D) \cdot 2}{n \cdot A_c \cdot f_s} \quad . \quad (3.17)$$

Equating (3.16) and (3.17) together and placing  $\xi$ , as a dummy variable yields:

$$\frac{V_{IN} \cdot D}{n \cdot A_c \cdot f_s} \cdot \mathbf{x} = \frac{(V_{IN} - V_c) \cdot (1 - D) \cdot 2}{n \cdot A_c \cdot f_s} \quad . \quad (3.18)$$

Solving for  $\xi$  will result in the flux density difference between the center leg and the two outside legs, seen in (3.19).

$$x = (1 - D) \cdot 2 \quad (3.19)$$

Therefore, the flux in the outside legs is  $2(1-D)$  times larger than the center leg.

### **3.5 Experimental Results for the Self-Driven Complementary-Controlled Full Bridge with Integrated Transformer**

The printed circuit board (PCB) design for any electronic circuit plays a critical role in the overall performance of a circuit. In the beginning design stages, most designers tend to overlook the importance of a well-planned PCB layout. The arrangement of the components on the PCB will greatly affect the circuit's performance, efficiency, and operation in a non-idealized and realistic world. In most compact, high-current, low-voltage DC/DC converters there are constraints on the converters footprint area, and the converter will have to fit into a given form factor. Two examples of a standardized form factor are a 1U voltage regulator module (VRM) that powers high-end microprocessors and a quarter brick that is a standard footprint for telecommunications industries server applications. This section presents the layout and experimental operation of a self-driven complementary-controlled full bridge circuit, with and without an integrated transformer in a quarter brick form factor.

Two self-driven complementary-controlled full bridge prototypes are built to test the performance of the integrated transformer. The design specifications for both boards are identical and each board was built to fit in the quarter brick form factor, shown in Fig. 3.21. The tested conditions for each board is  $V_{IN} = 36$  to  $75$  V with  $48$  V nominal,  $V_{OUT} = 1.2$  V,  $I_b = 70$  A at a  $750$  kHz switching frequency. The turns ratio on each power transformer is 12:1 and the two gate drive windings are 12:3.

The converters must operate up to  $75$  V input therefore, the primary switches chosen are Hitachi's HAT2175. Because the conduction losses of the secondary side

devices are dominant, five Hitachi HAT2165 synchronous rectifiers are used in parallel, each having a  $2.5\text{m}\Omega$  on resistance and  $15\text{ nC}$  gate charge. Each converter was built on a 9-layer 2-ounce copper printed circuit board. The output filter inductors for each board are  $200\text{ nH}$  and built with a Philips EI-18 core. The two inductors use an integrated magnetic technique. The windings for the cores are a single turn structure and only the two outside legs of the core has an air gap. The transformers on the discrete transformer board are as follows; the power transformer is Philips EI-14 3F4 material and the gate driver transformers are TDK EE9.5 PC50 material. The gate driver tiny MOSFET is Siliconix Si3900 and the gate drive resistor is  $0.5\ \Omega$ .

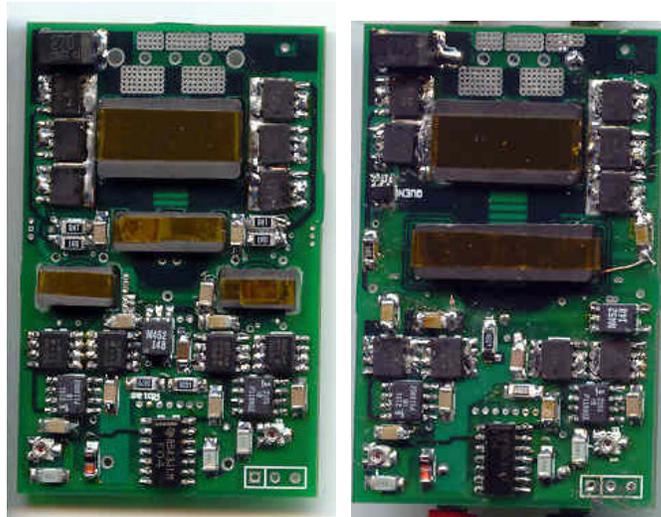


Fig. 3.21. Self-driven complementary-controlled full bridge, (left) discrete transformers and (right) integrated transformer.

The transformer on the integrated transformer is a custom designed core. The dimensions of the core are shown in Fig. 3.22. The operational waveforms for each board are shown in Fig. 3.23. Obviously the tested waveforms are nearly identical.

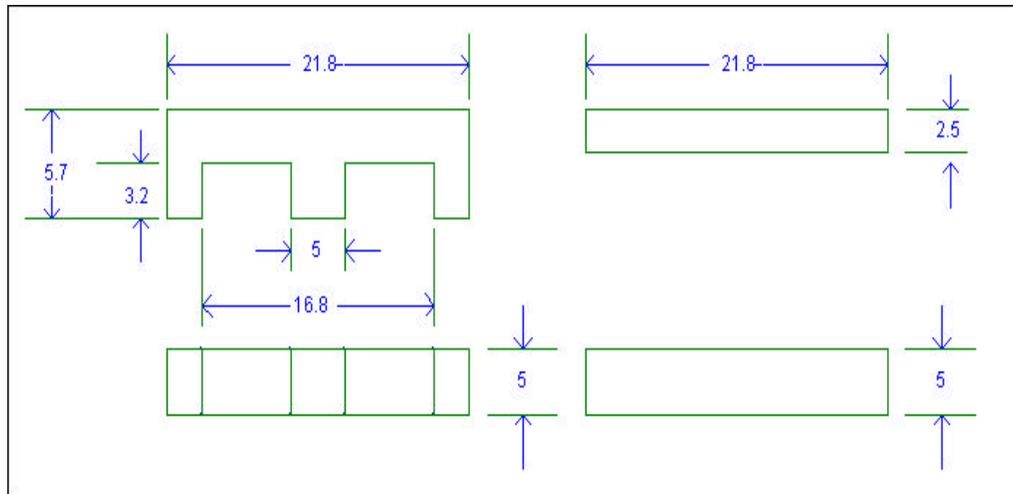
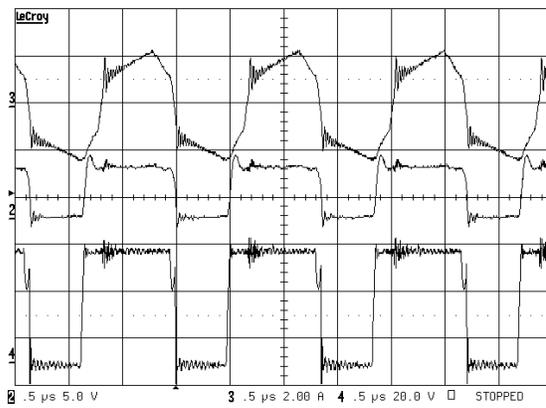
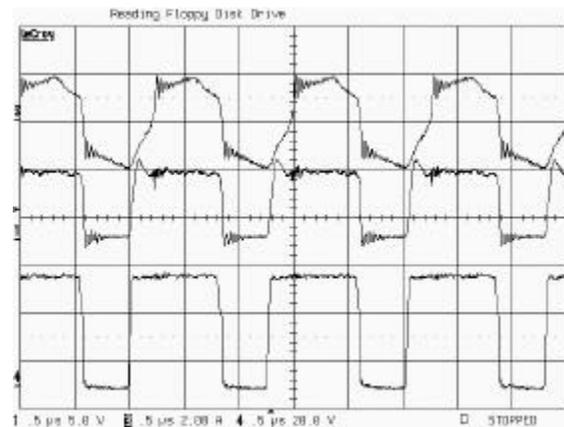


Fig. 3.22. Customized core for the integrated transformer design in (mm).

Referring to Fig. 3.23, the top waveform is the primary side current, the middle waveform is  $V_{GS}$  of the synchronous rectifiers and the bottom is the primary side Point A (B) corresponding to the gate drive signal. In each case the synchronous rectifier gate drive signal tracks the primary side.



(a)



(b)

Fig. 3.23. Operational waveforms. (a) The discrete transformer case, (b) Integrated transformer case. Top waveform is the primary side main transformer current, middle is the secondary side gate drive voltage, and bottom is Point A (B) corresponding to the gate drive waveform.

The efficiency curves for each prototype are shown in Fig. 3.24, and the loss breakdown shown in Fig. 3.25. As seen the efficiency of each converter is nearly identical and the overall efficiency at 57A load is 77%, which is less than adequate.

Therefore, a second iteration of the layout and design of the integrated transformer was performed to take advantage of the integrated transformer. The copper losses being the second highest loss, from Fig. 3.25, were targeted.

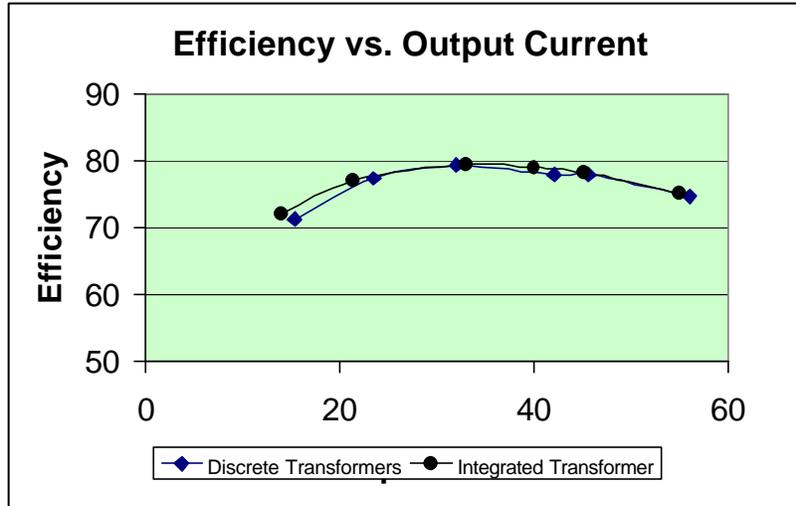


Fig. 3.24. Efficiency curve comparison.

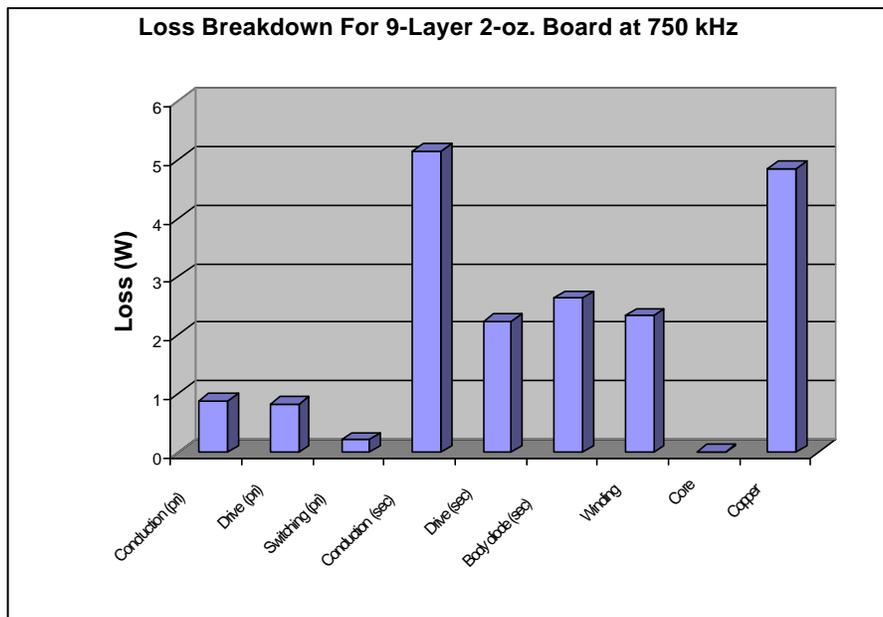


Fig. 3.25. Loss breakdown for the 9-Layer 2-oz. copper integrated transformer case.

To lower the copper losses, a 12-Layer 2-oz. copper board was used and the integrated transformer width was increase to 26 mm. The extra copper not only reduces the trace resistance but also acts as a heat sink. To lower the parasitic trace inductor in

the synchronous rectifier gate driver path, the gate drive signal is interleaved with the signal for the tiny MOSFET. The interleaving effect will lower the parasitic inductance. Finally, the output inductors were split into two discrete inductors to allow the sources of the synchronous rectifiers to nearly touch and be as close as possible to the ground pin. Having the sources of the synchronous rectifiers close to one another, and to the ground pin, lowers the overall high current ground path. The improved PCB layout is shown in Fig. 3.26, and the efficiency curve is shown in Fig. 2.27. Comparing the efficiency curve from Fig. 3.24 to Fig. 3.27, there is a 10% increase in efficiency. Seen in Fig. 3.27 the efficiency at 70A is greater than 86%.

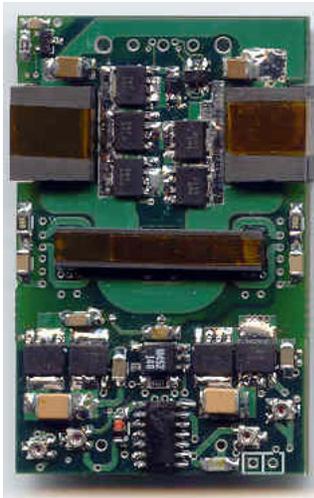


Fig. 3.26. Improved layout of the self-driven complementary-controlled full bridge with the integrated transformer.

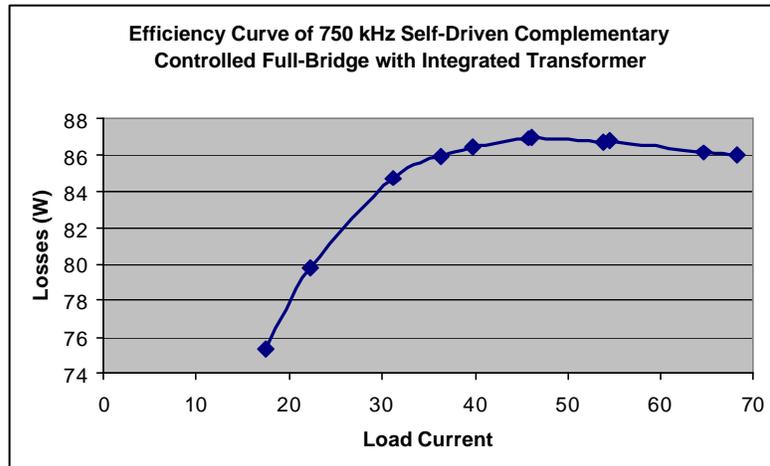


Fig. 3.27. Efficiency curve for the improved layout of the self-driven complementary-controlled full bridge with the integrated transformer.

### 3.6 Summary

This chapter presents and analyzes the concept of integrating two gate drive transformers and a power transformer into a single core structure. To lower the leakage inductance each passive secondary winding is coupled directly to a discrete active primary winding. The secondary winding can then interleave with the primary winding, resulting in very tight coupling. Integrating the two gate drive windings into the power transformer will reduce the synchronous rectifier gate driver length and the overall size of the power stage. In an isolated compact DC/DC converter, such as a quarter brick, the footprint of the power stage is limited. The space savings the integrated transformer gives can be used to optimize the layout.

From the experimental results presented above, the self-driven complementary-controlled full bridge with the integrated transformer can operate at 750 kHz with greater than 86% efficiency at 70 A load current in a quarter brick form factor. The self-driven complementary-controlled full bridge has greater than 50% reduction in footprint compared to the push-pull forward converter presented in the first chapter, while maintaining greater than 85% efficiency at 70 A.

# Chapter 4

## Conclusions and Future Work

With the arrival of the age of the high speed and power hungry microprocessors, the point of load converter (POL) has become a necessity. The power delivery architecture has changed from a centralized distribution box delivering the entire systems power to a distributed architecture where a common DC bus voltage is distributed and then further converted down at the point of load. Two common distributed bus voltages are 12 V for desktop computers and 48 V for telecommunications server applications. This thesis focuses on the following issues in the 48 V bus systems: secondary side gate drivers, self-driven circuits, integrated transformers and magnetics.

### 4.1 Summary

There are many different gate driver techniques used today. In high frequency applications, the losses due to driving many paralleled MOSFETs can be significant. A conventional gate driver is shown to have unacceptable loss levels, where resonant and self-driven gate drive circuits have loss saving abilities. In a compact, isolated power supply, a resonant gate driver can be bulky and difficult to drive, where self-driven gate drivers have small and simple structures. A new method of self-driven synchronous rectification is introduced in chapter 2. The self-driven circuit has gate drive loss saving properties, a fast driving speed, and can hold synchronous rectifiers on at the correct gate drive voltage, where the state-of-the-art self-driven gate drive circuit cannot. The self-driven gate driver loss savings are proportional to the gate resistance and parasitic trace inductance.

The self-driven circuit outlined in chapter 2 is used in conjunction with a complementary-controlled full bridge circuit with a current doubler secondary. The combination of the self-driven circuit and the complementary-controlled full bridge has a total of three transformers. Presented in chapter 3 is a novel method of integrating these three transformers onto a single core. The integrated transformer has three sets of windings that correspond to the main power transformer windings and the two gate driver windings. The integration method allows each set of winding to act as a discrete transformer. The integrated transformer saves valuable footprint area by eliminating two out of three magnetic cores, saves secondary side gate driver length and parasitic trace inductance.

## 4.2 Future Work

There are two interesting topics that need further development. These are an extension of the integrated transformer concept and a self-driven gate drive clamping method for wide input range converters.

### Extension of the integrated transformer

For a compact isolated design with a current doubler, the previous design proposed by Professor Wei Chen shown in Fig. 4.1 integrates the two output filter inductors with the main power transformer.

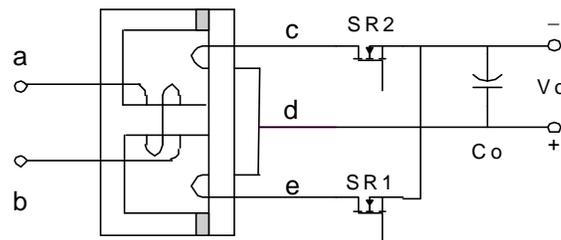


Fig. 4.1. Integrated current doubler rectifier proposed by Prof. W. Chen.

However, this structure is prone to larger leakage inductances because the secondary windings are not tightly coupled to the primary windings. Plus this winding structure requires the two outside legs of the EE/EI core to have an air gap and no air gap on the center leg. Gapping the two outside legs, and not the center, will cause a mechanical instability in the core, which makes it difficult to manufacture.

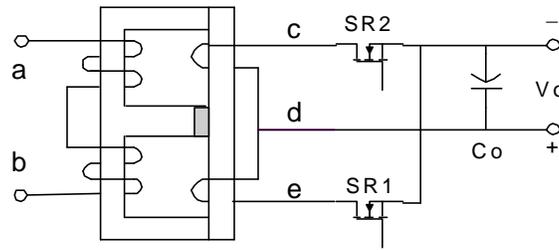


Fig. 4.2. Integrated current doubler rectifier proposed by Peng Xu.

Following Professor Wei Chen’s concept, Peng Xu then proposed splitting the primary windings to the outside legs of the transformer, seen in Fig. 4.2. The leakage inductor is reduced due to better coupling from primary to secondary windings. The core is gapped only in the center leg, which leads to two identical output inductors. Due to an interleaving effect, the AC flux in the center leg is reduced.

Incorporating Professor Wei Chen’s and Peng Xu concepts together, with the proposed integrated transformer, the following two converters can be realized, seen in Fig. 4.3 and Fig. 4.4.

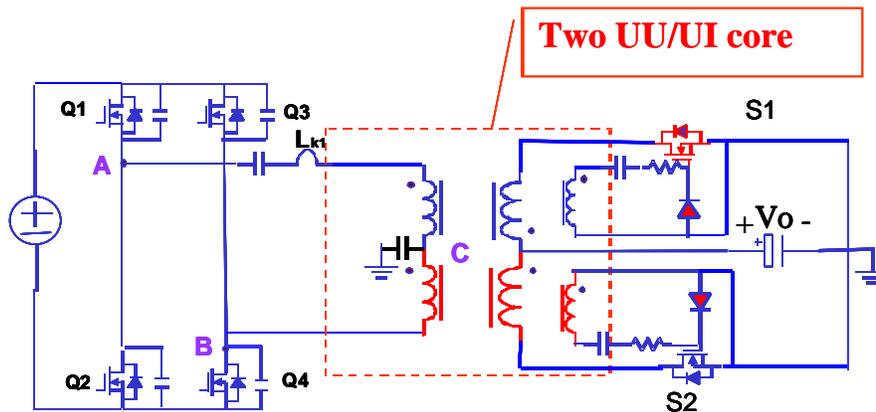


Fig. 4.3. Modified power stage of the self-driven complementary-controlled full bridge using two UU/UI cores to realize the two gate drive transformer, two output inductors, and the power transformer.

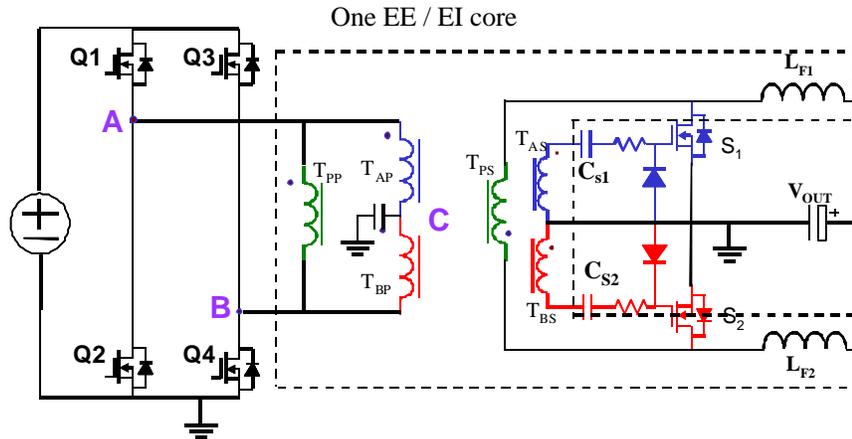


Fig. 4.4. Modified power stage of the self-driven complementary-controlled full bridge using a single EE/EI to realize the two gate drive transformer, two output inductors, and the power transformer.

The circuits in Fig. 4.3 and Fig. 4.4 have yet to be tested in hardware. Optimizing the winding structure and layout in a form factor such as a quarter brick or eighth brick would be a difficult, but worthwhile, task.

### A self-driven gate drive clamping method for wide input range converters

The self-driven gate drive voltage is dependant on the input voltage and the duty cycle. As the input voltage range varies, the voltage seen from the gate-to-source of the synchronous rectifiers also varies. For example, for a 36- to 75-volt input range on the self-driven complementary-controlled full bridge, with a turns ratio on the power transformer of 10:1 and gate drive transformers of 5:1 will have the following synchronous rectifier gate drive voltages, shown in Table 4.1.

<b>Table 4.1.</b> Power losses for given gate drive voltages using a conventional gate drive circuit. Circuit parameters are for a single HAT2165 operating at 750 kHz.			
<u>Input voltage (Volts)</u>	<u>V<sub>GS(SR)</sub> (Volts)</u>	<u>Q<sub>GS(SR)</sub> (nC)</u>	<u>P<sub>LOSS</sub> (Watts)</u>
36	7.2	50	0.27
48	9.6	70	0.504
75	15	100	1.125

The power losses under the 75-volt input case are greater than four times the losses for the 36-volt input case. The self-driven method illustrated in chapter 2 can save some of these gate drive losses, but the high line input case will always have more losses. Therefore, some method to reduce the high line input voltage synchronous rectifier gate drive losses are necessary.



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